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Design and test of a dynamic balancing machine for small rotors.

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DESIGN AND TEST
OF A DYNAMIC BALANCING MACHINE
FOR SMALL ROTORS

BY

JAMES E. HAMMERSTONE

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J. E. Hammerstone

DESIGN AND TEST

OF A DYNAMIC BALANCING MACHINE

FOR SMALL ROTORS

by
James E. Hammerstone,

Lieutenant, United States Navy

Submitted in partial fulfillment
of the requirements
for the degree of
MASTER OF SCIENCE

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the thesis requirements for the degree of

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United States Naval Postgraduate School

PREFACE

Volumous amounts of work have been done on the design of
sensitive dynamic balancing machines for rotors of various shapes
and sizes. The degree of sensitivity obtained greatly complicates
the design and adds to the cost of such a machine.

The purpose of this project is the design of a simple but
yet satisfactory balancing machine which will detect and locate
very small static and dynamic unbalances.

During the period from October, 1949 through April 1950,
the author designed a novel type balancing machine and carried out
tests to obtain experimental verification of the theoretical
analysis used. Calibration curves were also plotted for future
balancing problems in small rotors. This work was conducted
at the United States Naval Post Graduate School.

The project was suggested by Professor Ernest K. Gatcombe
and the author gratefully acknowledges his timely assistance and
guidance throughout the entire project. Acknowledgements are also
extended to the U. S. Naval Engineering Experiment Station for the
manufacture of the machine, for the use of film and galvanometer
units; to the David Taylor Model Basin for the use of Strain In-
dicators; to the Public Works Division of the U. S. Naval Academy
for making tracings; to Mr. J. A. Oktavec of the Mechanical Force
of the Post Graduate School for his workmanship in making special
apparatus, and to other members of the faculty for loan of instruments.

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LIST OF SYMBOLS

W - Weight

Z - Weight

W_1 - Weight

Z_1 - Weight

C_L - Weight

C_R - Weight

a - Distance

L - Correction plane

R - Correction plane

M - Mass

μ - Mass per unit length

ℓ - Length

Y_o - Static deflection

X - Coordinate

Y - Coordinate

E - Modulus of elasticity

I - Moment of inertia

$P.E$ - Potential energy

$K.E$ - Kinetic energy

ω_n - Natural angular frequency

m - $\mu \ell$

h - Distance

K_1 - Spring constant of beam #1

K_2 - Spring constant of beam #2

θ - Angle of rotation

g   - Acceleration of gravity

I_{cg} - Moment of inertia of the system about an axis in the xy plane through the center of gravity.

K_e - Equivalent spring constant for linear motion

K_{te} - Equivalent spring constant for torsional motion.

K_c - Coupling constant

f_1 - Natural frequency for linear vibration.

f_2 - Natural frequency for torsional vibration.

CHAPTER I

INTRODUCTION

It is impossible to manufacture a rotor in which the mass is so evenly distributed, that vibrations will not occur at the bearing supports. Therefore, to decrease bearing wear, failure due to vibration, etc., it becomes necessary to balance rotors both statically and dynamically.

A study of the engineering literature portrays many varieties of machines designed for static and dynamic balancing. They range from the comparatively simple machine of Laweczbeck Heyman, where the rotor is carried on spring supports, to the fine but complicated Gisholt Westinghouse machine which records the unbalance by a system of electrical networks.

In this day and age balancing requirements are so diverse that it is difficult to use one scheme for universal application. Therefore, it is not unusual to find that a new machine must be designed to accommodate a specific balancing problem.

The specific requirements for this machine is to accommodate light weight shafts ranging from two to four pounds in weight; seven to fifteen inches long and one quarter to one inch diameter.

Since the utmost requirement of a dynamic balancing machine is sensitivity, all effort throughout this project was funneled toward this requirement.

CHAPTER II

ELEMENTARY BALANCING CONSIDERATIONS

1. Nature of balancing.

   Balancing operations are of two types, static and dynamic.
Static balancing requires only the application of a single
correction weight, whereas dynamic balancing requires two
correction weights added in proper positions, in predetermined
planes at some distance apart and perpendicular to the axis
of rotation.

2. Static Unbalance.

   Static unbalance sometimes called force unbalance, when
of appreciable magnitude, is observed when the unbalanced part
is mounted on horizontal knife edges as in Figure 1.
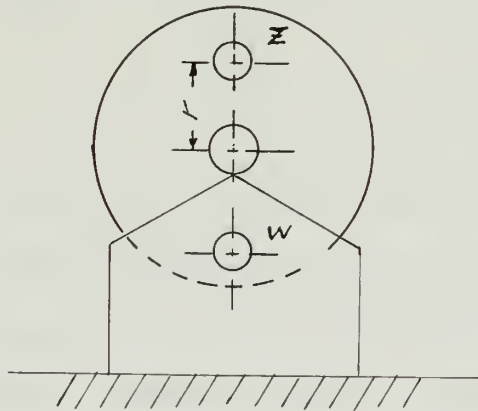
FIG 1

If the disc of Figure 1 has an unbalance equal to Wr $oz-in$;
a correction weight (Z) may be added 180° from W at such a
distance from the center that $Zr = Wr$.

3. Dynamic Unbalance.

The presence of dynamic unbalance or moment unbalance
occurs only when the piece is rotating. This type of unbalance
is caused by two equal weights at equal distances from the
rotational axis and located on opposite sides and opposite ends
of the rotating shaft as in Figure 2.
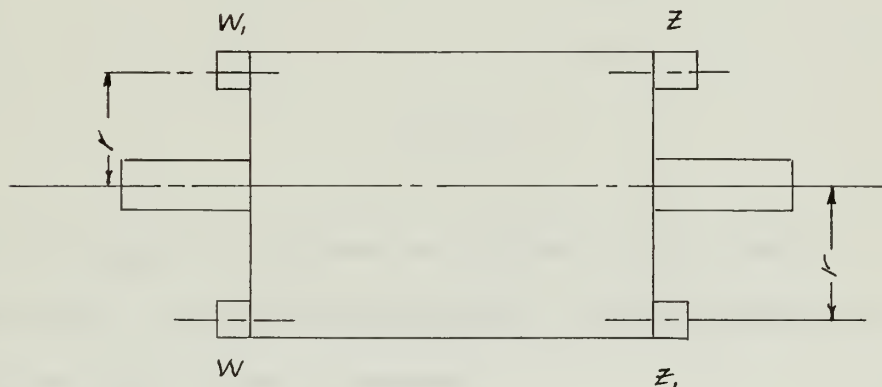
FIG 2

If in Figure 2, $Wr = Zr$ the piece will be in perfect
static balance, but on rotation, centrifugal forces due to
the weights will cause the ends of the cylinder to move in
opposite directions. The piece will be in dynamic unbalance.
By the addition of weights W_1 and Z_1 a compensating couple will
be introduced and the piece will be in dynamic balance.

4. Static and Dynamic Unbalance.

Rotating parts usually have both static and dynamic
unbalance. This is illustrated in Figure 3.

FIG 3

Let W in Figure 3 represent an unbalance in the rotor.
The planes L and R are representative planes in which cor-
rections can be added or removed.

If W is in the relative position as shown in Figure 3
two weights can be added in planes L and R to correct for both
static and dynamic unbalance. For static balance $C_L + C_R = W$.
For dynamic balance a moment equation of equilibrium can be
satisfied.

CHAPTER III

MATHEMATICAL ANALYSIS
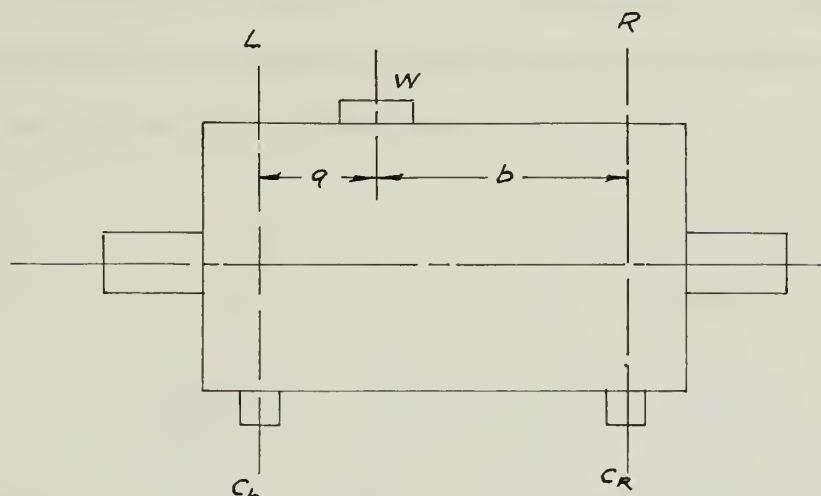
1. Problem.

The machine designed for this project can be simply
described as two rectangular cantilever beams supporting
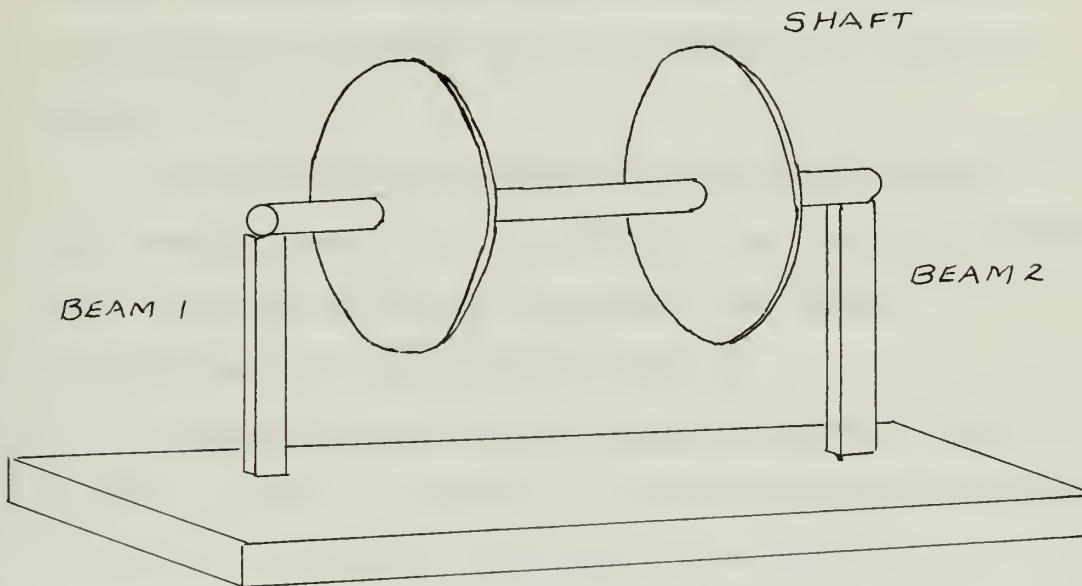a rotating shaft as illustrated in Figure 4.

FIG 4

If the shaft in Figure 4 is both statically and dynamically
unbalanced, the supporting beams 1 and 2 will experience a
vibratory motion when the shaft is rotated. This vibratory
motion will be extremely slight for minute unbalance, but if
the system is forced to vibrate at it's natural frequency a
maximum vibration, and therefore a max deflection of the beams

will occur.

However, to amplify the motion of the beams at resonant conditions of the system, requires that the natural frequency of the system be relatively high.

The mass of Beam#1, and Beam#2 must be considered when calculating the natural frequency of the system. To simplify matters, Rayleigh's method is used in finding what fraction of the beam's mass should be added to the concentrated mass supported at it's end.

The problem then is to determine the equivalent mass of each beam to be added to the concentrated mass they are supporting and to determine the natural frequency of the system.

2. Determination of the equivalent masses.*

In determining the equivalent masses, consider a cantilever beam EI of length (1) and mass ($\mathcal{Y}$) per unit length (total mass $m = \mathcal{Y}_1 \mathcal{l}$) carrying a concentrated load at its end as shown in Figure 5.
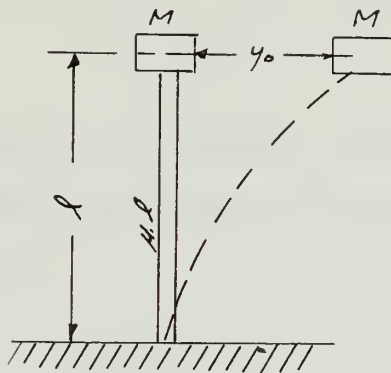

FIG 5

*See reference 3, 7.

Let mass (M) be deflected in the Y direction an amount y_0.
By Rayleigh's method (which is an approximation but sufficiently
accurate for this problem), it is assumed that the deflection
curve has the shape $y = y_0 \left(1 - \cos \frac{\pi x}{\ell}\right)$ imposing the following
end conditions:

$$x = 0, \quad y = 0$$

$$x = \ell \quad y = y_0$$

For the system the potential energy is

$$P.E. = \frac{EI}{2} \int_0^\ell \left(\frac{d^2 y}{dx^2}\right)^2 dx \qquad \text{\# (1)}$$

The kinetic energy is

$$K.E. = \frac{1}{2} y_1 \omega_n^2 \int_0^\ell y^2 dx + \frac{1}{2} M \omega_n^2 y_0^2 \qquad \text{\# (2)}$$

By substituting the assumed deflection curve then,

$$P.E. = \frac{EI y_0 \pi^4}{64 \ell^3} \quad \text{and} \qquad (3)$$

$$K.E. = \frac{1}{2} \omega_n^2 y_0^2 \left[M + m\left(\frac{3}{2} - \frac{4}{\pi}\right)\right] \qquad (4)$$

By the theory of Conservation of Energy, the Kinetic Energy =
Potential Energy, therefore equations (3) and (4) can be solved
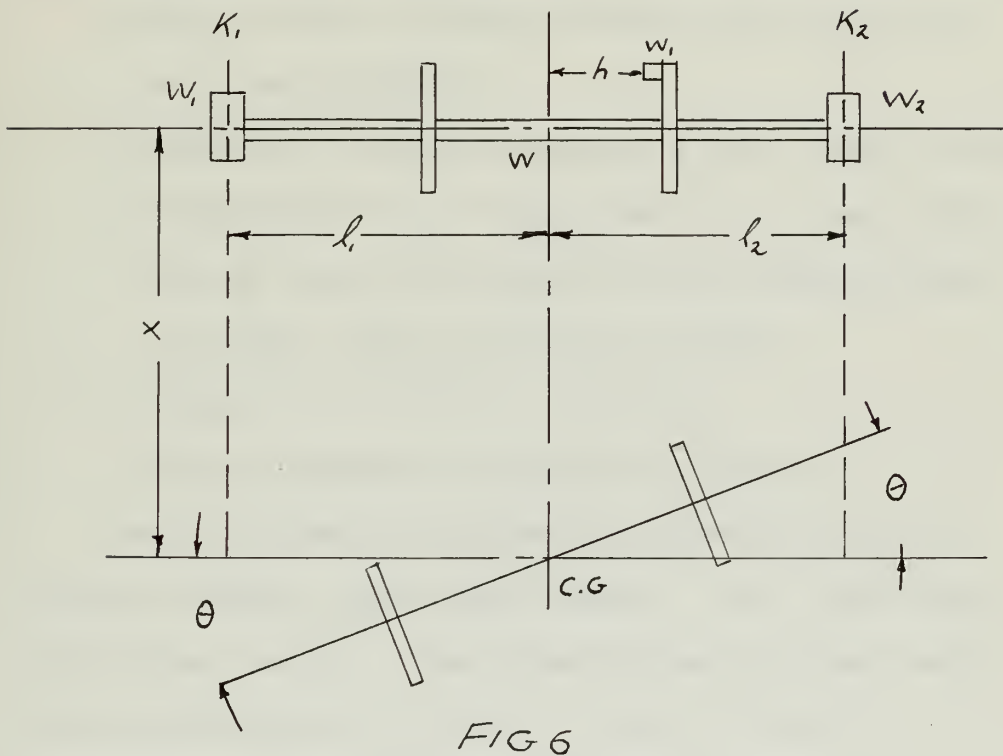for a natural frequency.

$$\omega_n^2 = \frac{3.03 \, EI}{\ell^3 (M + .23 m)}$$

See appendix

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Equation (5) shows that 23% of the mass $\mu_1 \ell$ of the beam should be added to the concentrated mass (M).

3. Equations of motion for the system*.

For setting up the equations of motion of the system a plan view of Figure 4 is utilized.

FIG 6

The system illustrated in Figure 6 , depending upon the constraints placed on it can have two degrees of motion. If the system is displaced it can have linear motion in the (X) direction, and a rotation about the center of gravity if the condition exists such that $K_1 \ell_1 \neq K_2 \ell_2$.

Note* See reference 5.

In deriving the equations of motion the following
assumptions are made:

1. The system is constrained to move in a linear
   direction along the (X) axis and a rotation in the
   xy plane about an axis through the center of gravity.

2. The rotor has sufficient stiffness that deflections
   in it are negligible.

3. Distances ℓ_1 and ℓ_2 from supports 1 and 2 remain constant.

4. The material is not stressed beyond its elastic limit
   and the beams do not deform in their built-in ends,
   thus their spring constants are linear.

5. Damping of the system is neglected.

6. Spring constants K_1 and K_2 are not equal.

When the system is displaced a distance x and rotated
through an angle θ, beam 1 will deflect $x + \theta \ell_1$ and beam 2
will deflect $x - \theta \ell_2$. The forces exerted by beams 1 and 2
will respectively be $(x + \theta \ell_1) K_1$ and $(x - \theta \ell_2) K_2$.
Therefore the equation of motion in the x direction is

$$\frac{W}{g} \frac{d^2 x}{dt^2} + x(K_1 + K_2) - (\ell_2 K_2 - \ell_1 K_1)\theta = 0 \tag{6}$$

By taking moments about the center of gravity the equation
of motion in the torsional direction results in

$$I_{c.g} \frac{d^2 \theta}{dt^2} + (K_2 \ell_2^2 + K_1 \ell_1^2)\theta - x(K_2 \ell_2 - K_1 \ell_1) = 0 \tag{7}$$

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In equations (6) and (7) the following symbols are defined.

W = total weight of the system

g = acceleration of gravity

x = linear displacement

θ = angular displacement

K_1= Beam 1 spring constant

K_2= Beam 2 spring constant

l_1= distance of Beam 1 from c.g.

l_2 = distance of Beam 2 from c.g.

$I_{c.g}$ = moment of inertia of the system about an axis in the xy plane through the center of gravity

~~c.g.~~ c.g. = center of gravity of the system

4. Solution for the natural frequency of the system.

A solution of equation (6) and (7) can be obtained in
the form

$$x = x_0 \sin \omega t$$

$$\theta = \theta_0 \sin \omega t$$

Making these substitutions result in:

$$- \frac{W}{g} x_0 \omega^2 + x_0 (K_1 + K_2) - \theta_0 (l_2 K_2 - l_1 K_1) = 0 \qquad (8)$$

and

$$- I_{c.g} \theta_0 \omega^2 + (K_2 l_2^2 + K_1 l_1^2) \theta_0 - (K_2 l_2 - K_1 l_1) x_0 = 0 \qquad (9)$$

It is seen that $K_1 + K_2$ will be the equivalent spring constant
for linear motion; for simplification let $K_e = K_1 + K_2$. The

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equivalent spring constant for torsional motion is

$$K_{te} = \left(K_2 \ell_2^2 + K_1 \ell_1^2 \right)$$

If $\left(K_2 \ell_2 - K_1 \ell_1 \right) = K_c$ where K_c is designated as the coupling constant, then equations (8) and (9) reduce to

$$\frac{W}{g} X_o \omega^2 - K e X_o + K_c \Theta_o = 0 \tag{10}$$

$$\frac{I}{c.g} \Theta_o \omega^2 - K_{te} \Theta_o + K_c X_o = 0 \tag{11}$$

Eliminating X_o and Θ_o results in the equation for the natural frequencies of the system.

$$\omega^4 - \left(\frac{K e g}{W} + \frac{K_{te}}{I_{c.g.}} \right) \omega^2 + \left(\frac{K e K_{te} - K c^2}{W I_{c.g.}} \right) g = 0 \tag{12}$$

When $K_2 \ell_2 = K_1 \ell_1$ the coupling constant is zero and equations (10) and (12) reduce to their uncoupled form. When the vibrations are not coupled each acts independently of the other. For this special case the following frequencies would result.

$$f_1 = \frac{1}{2\pi} \sqrt{\frac{K e g}{W}} \tag{13}$$

$$f_2 = \frac{1}{2\pi} \sqrt{\frac{K_{te}}{I_{c.g.}}} \tag{14}$$

In general the coupling will raise the higher frequency and lower the lowest frequency. If the coupling is small the effect on frequency is not very great.

5. Amplitude of vibration.

If in Figure 5 a small unbalanced weight (w_1) exists in the rotor and is located a distance (h) from the center of gravity, the system can be forced to vibrate at a natural frequency.

The equations for forced vibration neglecting damping are:

$$\frac{W}{g}\frac{d^2x}{dt^2} + x\left(K_1 + K_2\right) - \left(\ell_2 K_2 - \ell_1 K_1\right)\theta = \frac{W_1}{g} r\omega^2 \sin\omega t \tag{15}$$

$$I_{c.g.}\frac{d^2\theta}{dt^2} + \theta\left(K_2\ell_2^2 + K_1\ell_1^2\right) - x\left(K_2\ell_2 - K_1\ell_1\right) = \frac{W_1}{g} r\omega^2 h \sin\omega t \tag{16}$$

A solution can be in the form

$$x = A \sin\omega t$$

$$\theta = C \sin\omega t$$

Substituting the symbols for the spring constants, the constants A and C are determined as

$$A = \frac{\frac{W_1}{g} r\omega^2 \left[\left(K_{te} - I_{cg}\omega^2 + hK_c\right)\right]}{\left[\left(K_e - \frac{W}{g}\omega^2\right)\left(K_{te} - I_{cg}\omega^2\right) - K_c^2\right]}$$

$$C = \frac{\frac{W_1}{g} r\omega^2 \left[K_e + h\left(K_e - \frac{W}{g}\omega^2\right)\right]}{\left[\left(K_{te} - I_{cg}\omega^2\right)\left(K_e - \frac{W}{g}\omega^2\right) - K_c^2\right]}$$

Therefore

$$x = \frac{\frac{W_1}{g} r\omega^2 \left[K_{te} - I_{cg}\omega^2 + hK_c\right]}{\left[\left(K_e - \frac{W}{g}\omega^2\right)\left(K_{te} - I_{cg}\omega^2\right) - K_c^2\right]} \sin\omega t. \tag{17}$$

$$\theta = \frac{\frac{W_1}{g} r\omega^2 \left[K_c + h \left(K_e - \frac{W}{g}\omega^2 \right) \right]}{\left[(K_{te} - I_{c.g}\omega^2)(K_e - \frac{W}{g}\omega^2) - K_c^2 \right]} \sin\omega t \qquad (18)$$

Since Beam 1 deflects $x + \theta \ell_1$, then

$$x_1 = \frac{\frac{W_1}{g} r\omega^2 \sin\omega t \left[K_{te} - I_{c.g}\omega^2 + hK_c \right] + \ell_1 \left[K_e + h\left(K_e - \frac{W}{g}\omega^2 \right) \right]}{\left[(K_{te} - I_{c.g}\omega^2)(K_e - \frac{W}{g}\omega^2) - K_c^2 \right]} \qquad (19)$$

Beam 2 deflects $x - \theta \ell_2$, then

$$x_2 = \frac{\frac{W_1}{g} r\omega^2 \sin\omega t \left[K_{te} - I_{c.g}\omega^2 + hK_c \right] - \ell_2 \left[K_c + h \left(K_e - \frac{W}{g}\omega^2 \right) \right]}{\left[(K_{te} - I_{c.g}\omega^2)(K_e - \frac{W}{g}\omega^2) - K_c^2 \right]} \qquad (20)$$

The deflections may be expressed as ratios:

$$\frac{x_2}{x_1} = \frac{\left[K_{te} - I_{c.g}\omega^2 + hK_c \right] - \ell_2 \left[K_c + h \left(K_e - \frac{W}{g}\omega^2 \right) \right]}{\left[K_{te} - I_{c.g}\omega^2 + hK_c \right] + \ell_1 \left[K_c + h \left(K_e - \frac{W}{g}\omega^2 \right) \right]} \qquad (21)$$

CHAPTER IV

APPROACH TO THE DESIGN PROBLEM

In designing a balancer capable of detecting and locating very small unbalances, certain important factors and problems which do not exist in more common balancing operations must be considered.

1. Problems in Design.

The problems encountered are:

(1) Designing into the machine a very high degree of sensitivity.

(2) Spinning the rotor in the balancing fixture. Mechanical couplings or belts introduce an unbalanced moment in the shaft after it has been removed from the balancer. Electrical drive is some times impossible or impracticable.

(3) Proper selection of bearings, such that they absorb the minimum of energy.

(4) Selecting a material for vibrating members which will withstand cyclic stresses without failure and have inherently low internal damping capacity.

2. Sensitivity of the machine.

Sensitivity may be defined as the amplitude of vibration of the system per unit unbalanced mass at unit radius.

Probably the most important factor affecting sensitivity is the overall damping. This presents a most difficult

situation and can hardly be solved without recourse to preliminary investigation. This investigation was beyond the scope of the project.

The overall damping is due to a variety of causes such as (a) air resistance (b) damping of supports (c) damping due to bearings (d) internal damping in the metal of the elastic members (e) energy loss in the foundation. Maximum sensitivity depends on reducing these to a minimum.

Damping (a) can usually be considered small, nevertheless, it still is impracticable to remove. In the case of (b) supports must be firmly fixed and fits between mating surfaces must be extremely close. In the case of (c) ball bearings should be avoided if possible, or very high precision instruments bearings with small clearances and minimum friction should be used. Perhaps the most important aspects are those of (d) and (e).

The main object in designing the suspension system is the choice of a material of low internal damping capacity. Although there are volumes of material published concerning the damping capacity of metals, knowledge in this sphere is still limited. Following is an exerpt taken from a letter received by the author from the Aluminum Company of America.

"From work that has been done on the subject it is generally accepted that aluminum alloys have extremely low damping capacities. Also it is definitely indicated that the damping capacity of most structures or assemblies is much more a function of the joints and connections than of the internal damping capacity of the material."

The main value of the published data is in showing
that certain materials, mainly brasses, bronzes, duraluminum
and hardened steel have low damping capacities, while ductile
materials, cast metals and plastics exhibit higher damping
characteristics..

The base upon which the suspension system is attached
should be of large mass and high stiffness so that its
natural frequency of vibration is well above that of the sus-
pension.

CHAPTER V

APPARATUS

Detailed plans of the balancing machine designed by the author are enclosed, attached to the back cover in Folder #1.

1. Suspension System.

A relatively high natural frequency of the system dictated the use of a light material for the vibrating members. The suspension system consists essentially of two cantilever beams extending vertically from the base. The beams, with a cross-section of 1/2" x 3/4", were constructed from duraluminum (*24ST-4*). They extend into the base with very close snug fits and anchored with nuts.

It was necessary to machine the beams to very close tolerances so their spring constants would be similar. To have them exactly equal, constitutes a difficult problem. The spring constant is affected not only by the dimensions but also by the non-honogeniety of the heat treatment, and differences in the fits of mating parts.

2. Power System.

Since mechanical couplings or belts used for spinning the rotor in the balancing fixture introduce unbalanced moments, small air driven wheels were used. Each end of the rotor journal rests between two such wheels, thus eliminating any unbalanced moment. The wheels were not machined integrally

due to problems arising in milling out the small buckets
which are spaced at 7° intervals around the periphery. The
buckets were milled in a solid aluminum disk and two flanges
were riveted to the disk for complete assembly. The method
is illustrated in Figure 7.
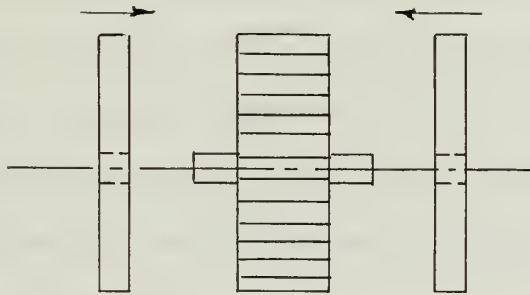
FIG 7

- The driving wheels are supported on R-2 Inch Series New
Departure ball bearings (1/8" I.D; 3/8" O.D.) They are driven
simultaneously by compressed air which follows a path through
a 1/8" drilled hole in the cantilever beams and cross supports.
It is then directed through 3/32" drilled holes in the wheel
support brackets leading to 3/32" diameter nozzles which project
air into the buckets at a 30° angle with the vertical.
3. Base.

The base of the machine was designed exceedingly heavy to
eliminate the necessity of bolting it down. A heavy base
is also advantageous in that its natural frequency is well above

that of the suspension system. Drilled and reamed holes
were provided to accommodate rotors of various length.

4. Test Shaft.

A test shaft was designed for calibrating the machine in
both static and dynamic unbalance. Two disks which provide
correction planes, in which known weights can be added, where
welded on the journal so they would be in close proximity to
the supports to produce maximum vibrations. Holes 1/8" diameter
where drilled and tapped in the face of the disks; spaced every
10° to accommodate small aluminum correction weights.

5. Protection device.

Since the rotor is free to move vertically upward off it's
driving wheels if the centrifugal force due to unbalances is
sufficient, it was necessary to provide for protection. The
protecting guard consists of two flat steel bars extending across
the rotor and supported at the base. The height of the guard
can be adjusted to accommodate different size rotors.

6. Vibration Pickup System.

From the theory of Strength of Materials a cantilever
beam with a concentrated mass on its end will have its maximum
flexural stress and therefore maximum strain at the built-in
end. Using this fact as a basis a system was set up to measure
the strains produced in the beams by the unbalanced rotor.

A relatively new, but very proficient system is the use
of wire strain gages. Figure 8 illustrates the type of gage
used.

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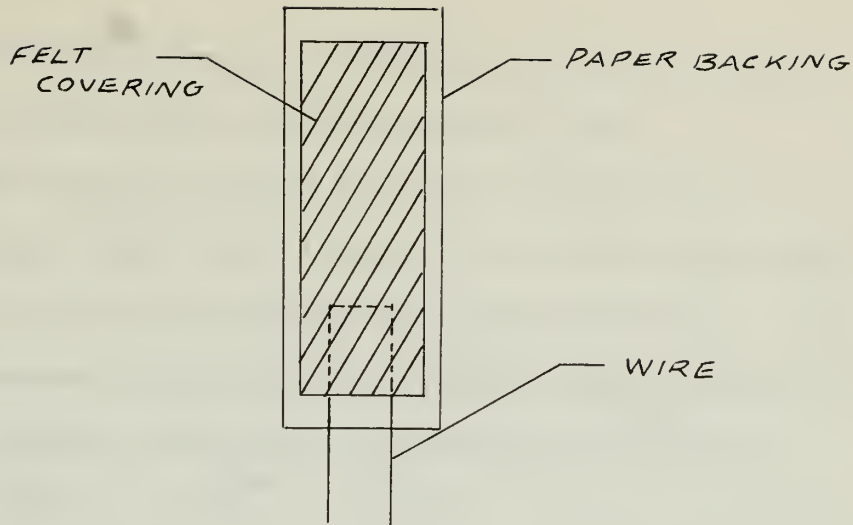
WIRE

FIG 8

The gage is of the A-1 SR-4 type manufactured by
Baldwin Locomotive Works, Philadelphia, Pa.

The fine wire forming the strain sensitive element is
usually constructed of .001" dia wire of some material having a
low thermal coefficient of resistivity. Alloys of the constantin
type are frequently used. The sensitive wire is covered with
a square piece of felt and supported by thin paper having good
mechanical and insulating characteristics.

The gage operates on the principle that a change in length
of sensitive wire produces a proportional change in its
electrical resistance. This fact may be expressed by

$$\frac{\Delta R}{R} = K \frac{\Delta \ell}{\ell}$$

Where R is the electrical resistance of the gage

Δr is the change in R due to strain

l is the length of gage

Δl is the change of length

K is a constant of proportionality.

-20-

To make an accurate measurement of the strain in the
beams a carrier type strain indicator was used.

The TMB type IA strain indicator developed by the
David Taylor Model Basin, utilizes two A-1 SR-4 strain gages with
a gage resistance of 120 ohms and a gage factor K of 2. They
are connected as adjacent arms in an A.C. bridge circuit; the
block diagram of the system is illustrated in Figure 9.

BRIDGE AND CALIBRATION
CIRCUIT
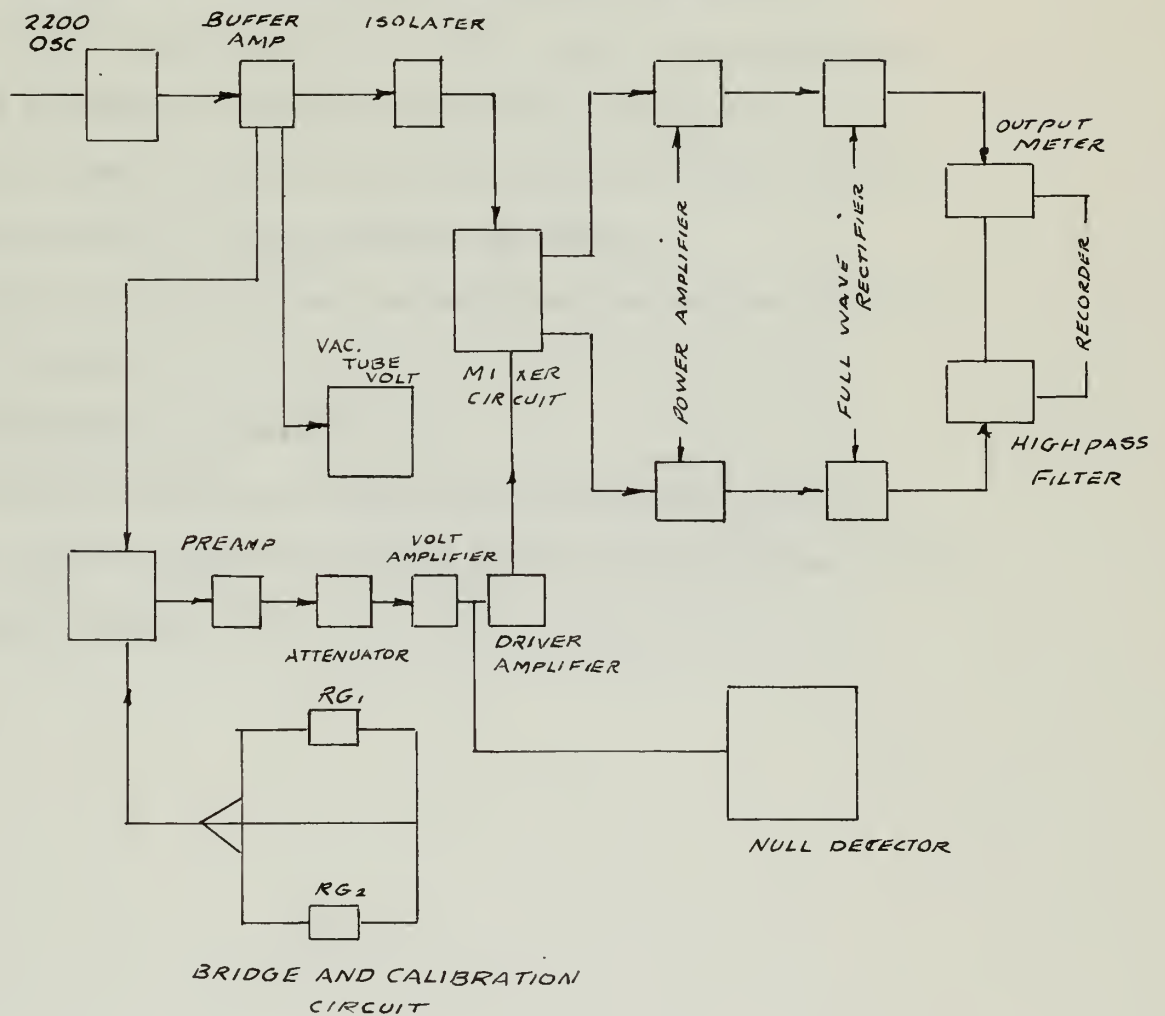
FIG 9

The indicator can be calibrated to read from 40 micro inches/inch
to 2500 micro inches per inch strain with an accuracy of$\pm$ .25%

For detail discussion on design and operation see reference
#2 in Bibliography.

7. Recorder.

The output of the strain indicator was fed into a
Universal Type P.A. Westinghouse Oscillograph. Since the
max output of the indicator is 4 mill. amps, a super-sensitive
string galvanometer sensitive to 23 mill. amps was used. Only
one such element could be located in the area, therefore simul-
taneous readings for each beam was impossible.

When two gages are used on each beam, the recorder indicates
double strain.

8. Arrangement of Apparatus.

Figure 10 shows an overall view of apparatus. Figure 11
gives a close up view of balancing machine. The recording machine
is shown in Figure 12.

FIG 10
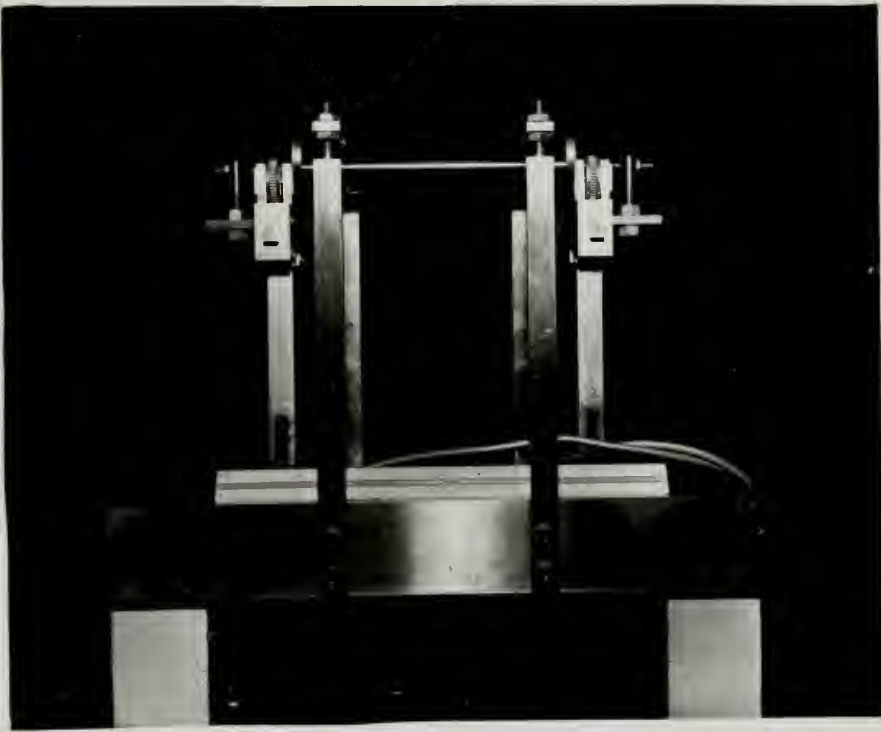
FiG 11.
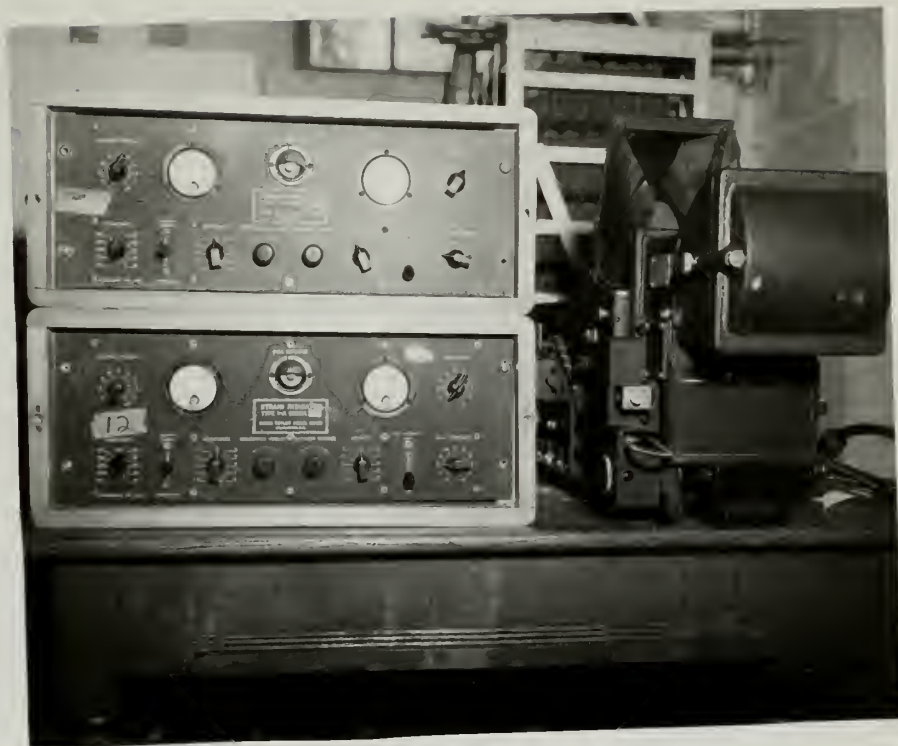
Fig 12

CHAPTER VI

PROCEDURE AND RESULTS

1. Preliminary tests.

Before running any tests on the machine, it was necessary to determine the mass of each part and their centers of gravity. From this data the moments of inertia were calculated. The spring constants of the beams were determined experimentally. The test shaft, not in perfect balance when manufactured, was dynamically balanced as best could be, considering that time was very limited.

2. Masses and centers of gravity.

Each component of the machine was weighed to the nearest 1/100 of a gram. The weights are listed in Table I, opposite to a sketch of each part.

Due to the complicity of the design, it was necessary to determine the centers of gravity by an experimental method. This consisted of finding the balance point of the parts or assembly of parts, by placing them on a knife edge which rested on a level plane. The centers were measured to the nearest 1/100".
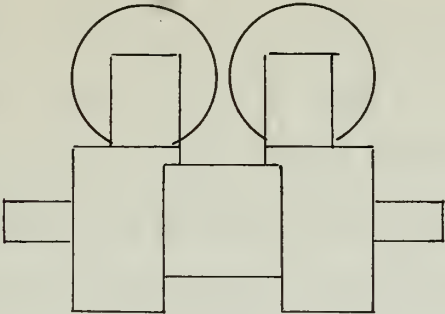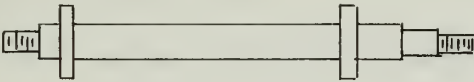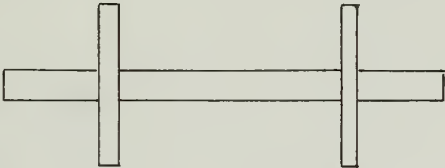
PART SKETCH	WEIGHT- GRAMS
CROSS SUPPORT and WHEEL ASSEMBLY 1.	296.90
CROSS SUPPORT AND WHEEL ASSEMBly 2.	296.70
BEAM 1	173.43
TEST ROTOR	1252.10
BEAM 2	173.35

TABLE I

3. Spring constants.

An attempt was made to measure the static deflection of the beams by an optical telescope. This proved unsuccessful since the deflections could be readonly to the nearest .0005".

An Ames-Jeweled Bearing Dial Indicator Gage which recorded
to the nearest 1/10,000 inch deflection did however prove
successful. Before making the measurements the machine was
assembled and aligned, then turned on its edge and clamped
to a drill press table. The beams were loaded statically such
that the line of force of the load passed through the center
line of the shaft. Tables II and III summarize data recorded
for each beam. Figures 13 and 14 show the graphs of deflection
plotted against load.

The spring constants determined:

$$K_1 = 335 \ lb/in$$

$$K_2 = 316 \ lb/in$$

Load (grams)	Deflection (inches)
50	.00025
100	.00060
200	.00125
500	.00325

TABLE II
Data for Beam 1

Load (grams)	Deflection (inches)
50	.00030
100	.00075
200	.00140
500	.00350

TABLE III
Data for Beam 2

SPRING CONSTANT CURVE
FOR
BEAM I

FIG 13
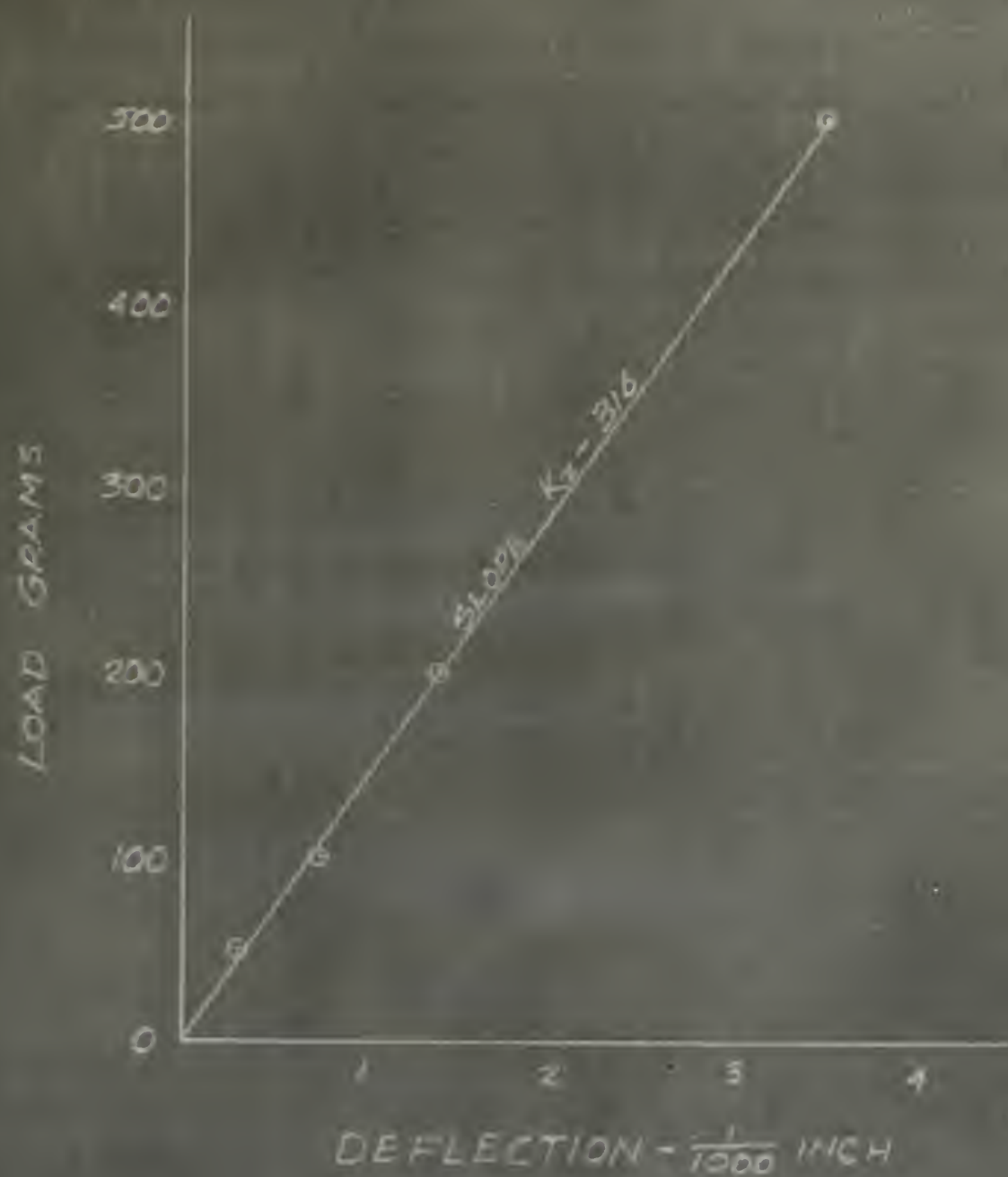
SPRING CONSTANT CURVE

FOR

BEAM 2

FIG 14

4. Natural Frequency.

The natural frequencies of the system were determined
experimentally by spinning the rotor to a speed above the
resonant condition and then letting it decelerate slowly.
When the amplitude of vibration reached a maximum as recorded
on the oscillograph, the speed was noted on the strobotac.
The system, as the theory predicted, experienced two natural
frequencies. The lowest at a speed of 2410 R.P.M. was due to
linear vibration and the highest at speed of 3020 R.P.M. was
due to torsional vibration.

Theoretically, the frequencies were calculated at
$f_1 = 43$ cps or 2580 rpm and $f_2 = 51.7$ cps or 3100 rpm.

These frequencies can be calculated in the following
manner using the theory in Chapter II.

The total weight of vibrating beam is .2456 lbs. and its
center of gravity is 3.4 inches from the base as shown in Figure 15.
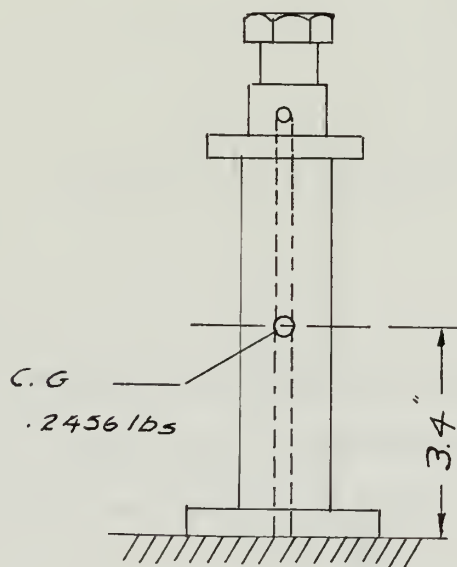
FIG 15

The total weight supported by each beam, disregarding the weight
of the shaft, is .870 lbs. This includes weight of wheels, supports
cross support and end bearings and holders. The center of gravity
of the beam assembly including the concentrated mass is 6.74"
from the base as shown in Figure 16.
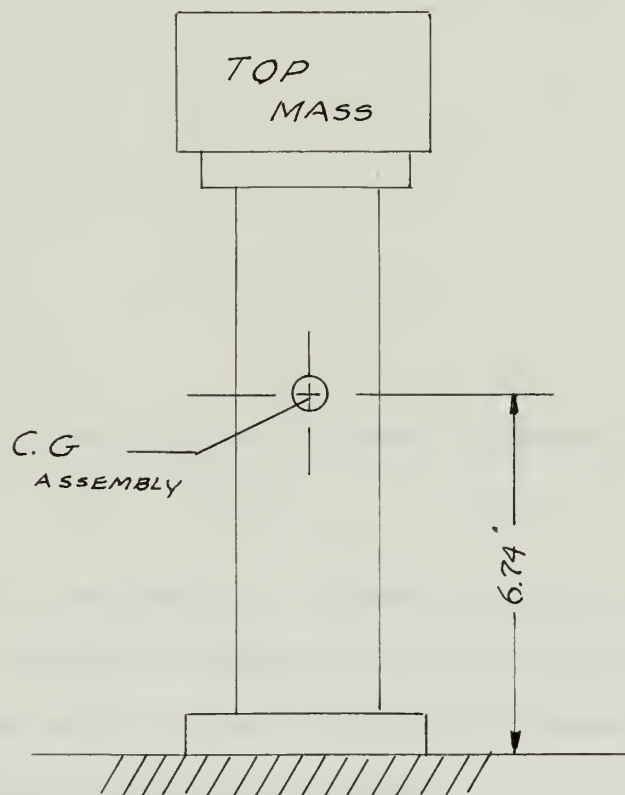
FIG 16

According to Timoshenko and DenHartog, as shown in equation 5,
Chapter II, the system can be considered as a cantilever beam
supporting a total load consisting of the concentrated load

plus 23% of the weight of the beam.

These considerations reduced the original system to
an equivalent consisting of a mass less beam of length 6.74"
supporting a load of .926 lbs. This is represented in Figure 17.

$$M + .23m = \frac{.926}{386}$$

6.74'
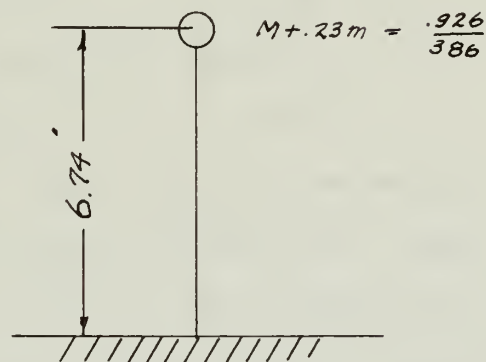
FIG 17.

If the above system is displaced its natural frequency is:

$$\omega_{n_1}^2 = \frac{3.03 \, EI}{\ell^3 (M + .23m)} \qquad \left(equation \ 5 \ Chapt. \ \overline{\mathbb{I}} \right)$$

where E and I are considered constant over length (ℓ).

Now, consider the weight of the shaft supported by the
beams. This weight is evenly distributed between the beams and
is concentrated 9.375" above the base.

Therefore, the mass $\left(M + .23m \right)$ must be transferred to this
point such that the resulting system will vibrate at the same
frequency as that shown in Figure 17. This transfer can be
illustrated as shown in Figure 18.

$$M_2 = f(M + .23m)$$

$$M_1 = M + .23m$$

$$l_2 = 9.375''$$

$$l = 6.74''$$

F, G 18
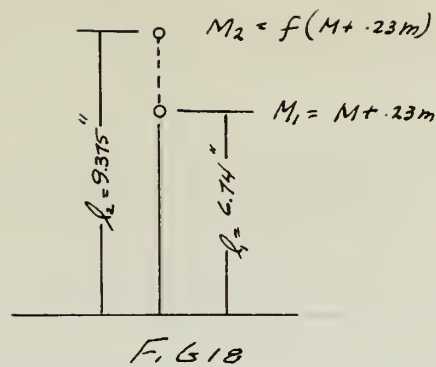
By using the same equation as before, then

$$\omega_{h_2}^2 = \frac{3.03\, EI}{l_2^3\,(M_2)}$$

Since $\omega_{h_2}^2 = \omega_{n_1}^2$ then

$$\frac{3.03\, EI}{l_2^3\, M_2} = \frac{3.03\, EI}{l_1^3\, M_1} \qquad \text{or} \qquad \frac{M_2}{M_1} = \frac{l_1^3}{l_2^3} \quad ; \quad M_2 = \frac{l_1^3}{l_2^3}\, M_1$$

Substituting the appropriate values.

$$M_2 = \left(\frac{6.74}{9.375}\right)^3 (.926) = .344\ lbs.$$

Therefore .344 lbs. must be added to the weight of the shaft or

the total weight supported by each beam = .344+1.385 = 1.729 lbs.

Before calculating the natural frequencies of the system

by equation 12, Chapter II, the moment of inertia of the system

about its center of gravity must be determined. Let Figure 19

represent the shaft and the point concentrated masses supported

by the beams.

DISK 1   DISK 2

M_1 C.G$_s$ M_2

l_1 l_2

K_1 K_2
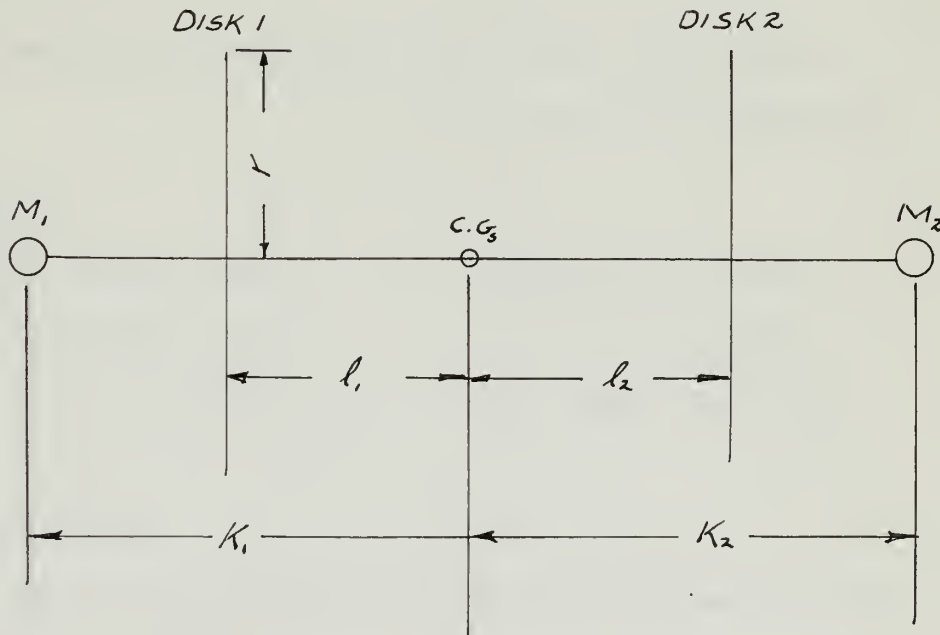
FIG 19

For disks $I_{c.g_s} = 2\left[M_{DISK}\dfrac{r^2}{4} + M_{disk}\,C_1^2\right] = .070\ units$

For $M_1 + M_2$ $\quad I_{M_1+M_2} = 2\left(M_1\ell_1^2\right) = .0304\ units$

Total $I_{c.g_s} = .104\ units$

Total weights of system = 3.458 lbs.

Using the following equation to find ω_n

$$\omega_n^4 - \left(\frac{Ke\,g}{Wr} + \frac{Kte}{I_{c.g.}}\right)\omega_n^2 + g\left(\frac{Ke\,Kte - Kc^2}{Wr\,I_{c.g}}\right) = 0$$

Substituting the following values:

$Ke = 316 + 335 = 651$

$Kc = (335 - 316)(4126) = 7835$

$Kte = (651)(4.126)^2 = 11,100$

$W_t = 3.458$

$I_{c.g} = .104$

$g = 386$

The resulting frequencies are f_1 (linear) = 2580 c.p.m.

$\qquad\qquad\qquad\qquad\qquad f_2$ (TORSIONAL) = 3100 c.p.m.

Summary.

Calculated values	Experimental values.
f_1 = 2580 or 43 cps	f_1 = 2410 or 40 cps
f_2 = 3100 or 51.7 cps	f_2 = 3020 or 50.3 cps

5. Calibration.

Before running calibration tests the rotor was balanced as nearly perfect as possible. This was a trial and error method of adding very small weights of putty in both disks.

To illustrate the method of calibration refer to Figure 20.
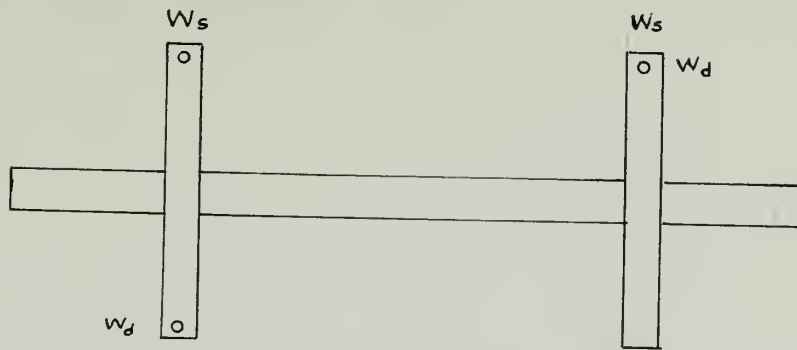
$F_I G_{20}$

Static calibration was conducted by adding equal weights
W_s as shown in Figure 20 at the same angle in each disk. The
rotor was decelerated very slowly through the natural frequency.
The high amplitude of vibration which exists near resonance
is the result of the energy fed into the system by the unbalanced
force, and obviously depends on the time taken in passing through
the critical range. Therefore, critical speed control is necessary,
Photographs were taken when the amplitude was maximum. The
oscillographs for each beam are shown on the following pages.
To obtain the data for the calibration curves, the zero un-
balance condition was used as a base. The amount of strain
produced by each unbalance weight was determined by taking the
differences in amplitudes. Since double strain was recorded
the results were halved. Results are shown in Table IV and V.
Figures 21 and 22 are the calibration curves.

STATIC CALIBRATION

BEAM 1

FIG 21

STATIC CALIBRATION

BEAM 2

FIG 22

Unbalance (gram-cm)	Strain (micro inches/inch)
0	0
.594	9.00
1.005	12.00
1.395	14.25
1.800	21.00

TABLE IV
Beam #1 - Static

Unbalance (gram-cm)	Strain (micro inches/inches)
0	0
.594	7.5
1.005	12.0
1.395	18.0
1.800	24.5

TABLE V
Beam #2 - Static

Dynamic calibration was conducted in the same manner as described for static, except two equal weights were added, one in each disk 180° out of phase as shown in Figure 20. Results are shown in Tables VI and VII. Figures 23 and 24 are the calibration curves.

Unbalance (gram-cm)	Strain (micro inches/inch
0	0
.594	6.75
1.005	11.25
1.395	16.50
1.800	19.50

TABLE VI
Beam 1 - Dynamic
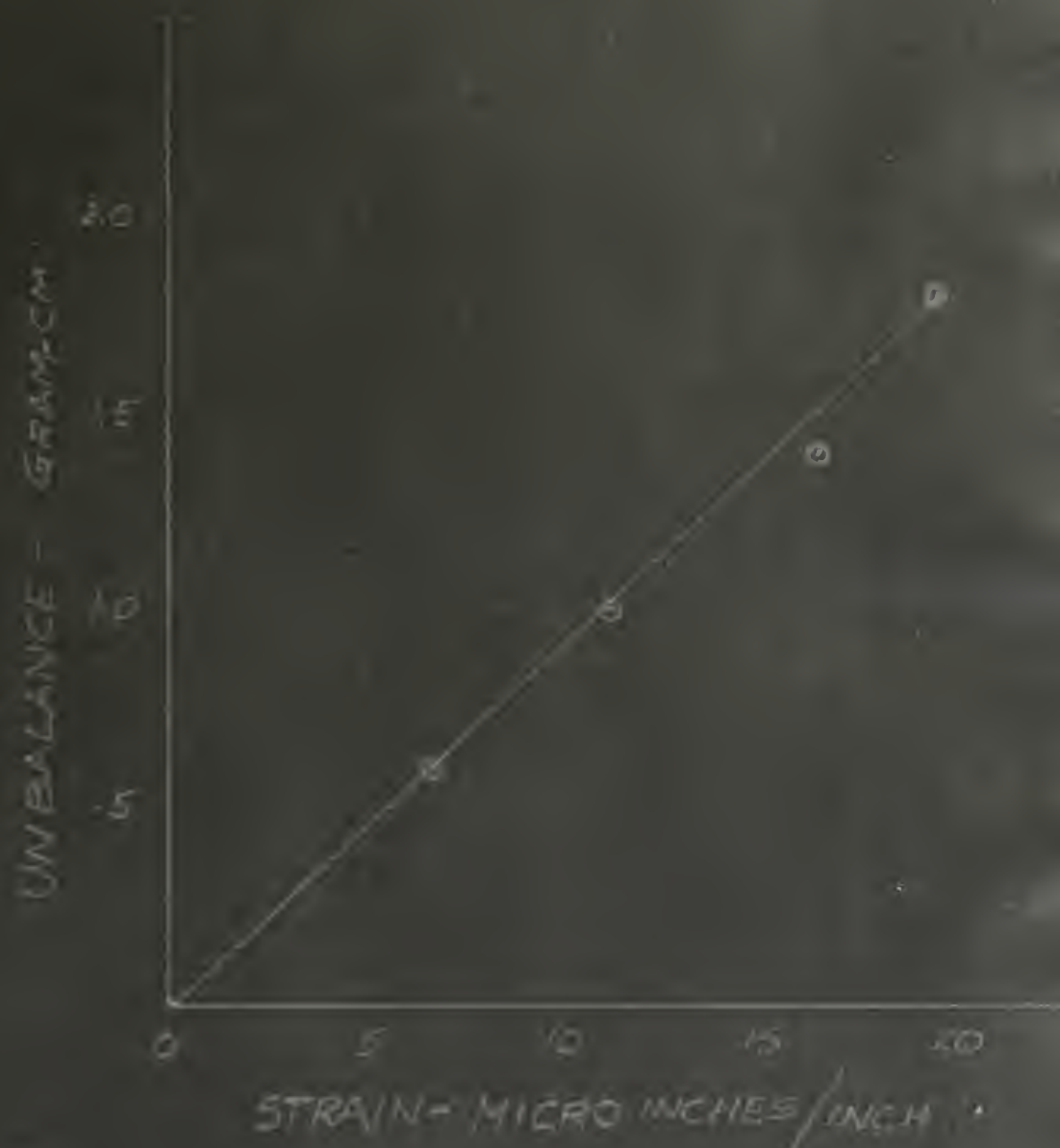
DYNAMIC CALIBRATION

BEAM 1

FIG 23
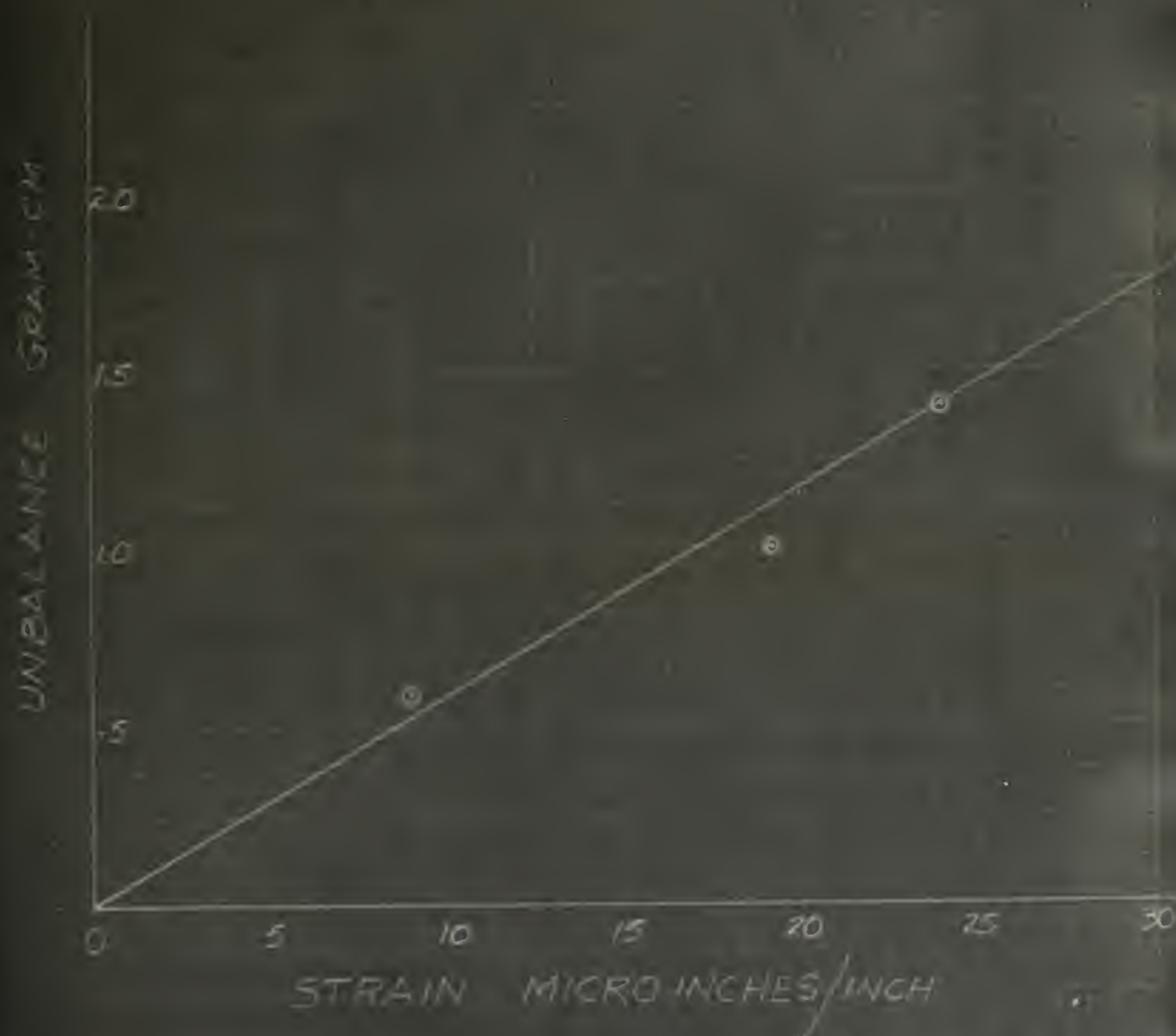
DYNAMIC CALIBRATION

BEAM 2

FIG 24

-71-

Unbalance (gram-cm)	Strain(micro inches/inch)
0	0
.594	9
1.005	19.25
1.395	24.0
1.800	32.25

TABLE VII
Beam 2 - Dynamic

For calibration of the machine, the wire strain gages were
placed as near as possible to the built-in ends of the beams.
The beams were scribed marked for future installation of gages.

6. Angular location of unbalance.

For very small unbalances it was necessary to use an optical
telescope focused on a hairline etched on the beams, to facilitate
seeing the movements of the system. When the disk and hairline
are stroboscopically illuminated, both the disk and hairline
can be made to appear stopped. If the hairline is stopped in
one of its maximum positions, the phase angle of the unbalanced
weight can be spotted. A telescope was not necessary for the
larger unbalanced weights. The unbalance was spotted at the end
of its horizontal radius when the beam was in its maximum
horizontal position which is in accord with the theory presented
by *Den Hartog.* .
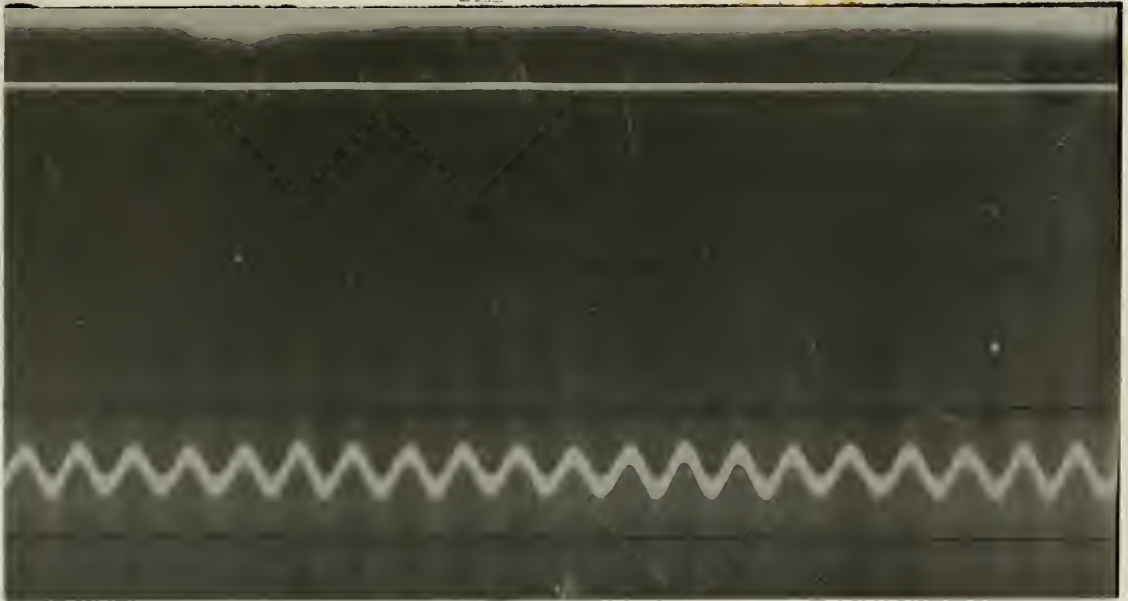
BEAM 1
O UNBALANCE

BEAM 2

O UNBALANCE.
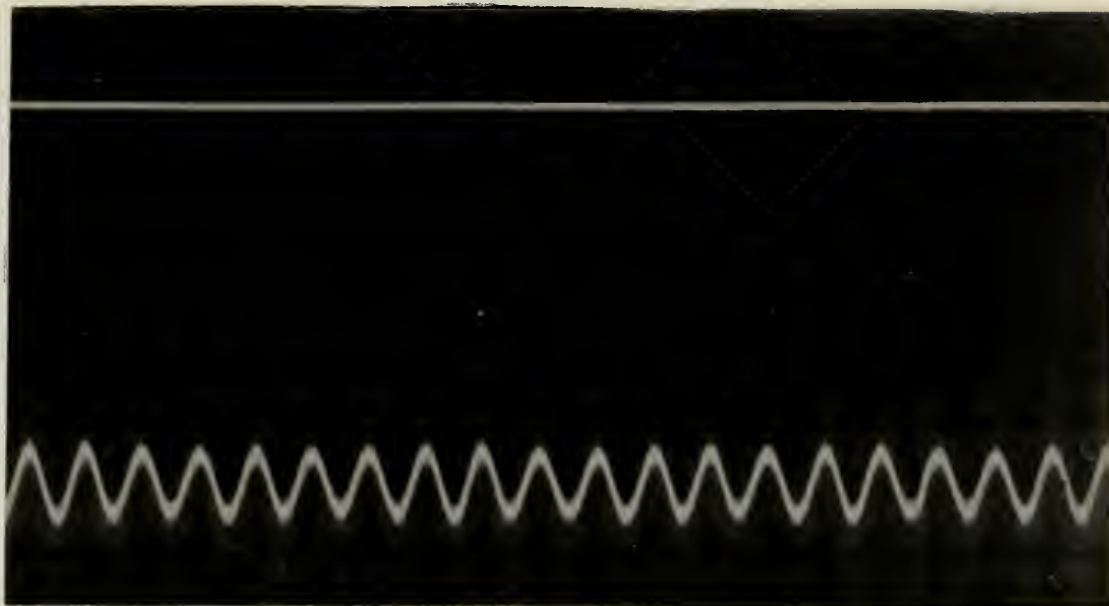
BEAM 1- STATIC
· 156 GRAMS.
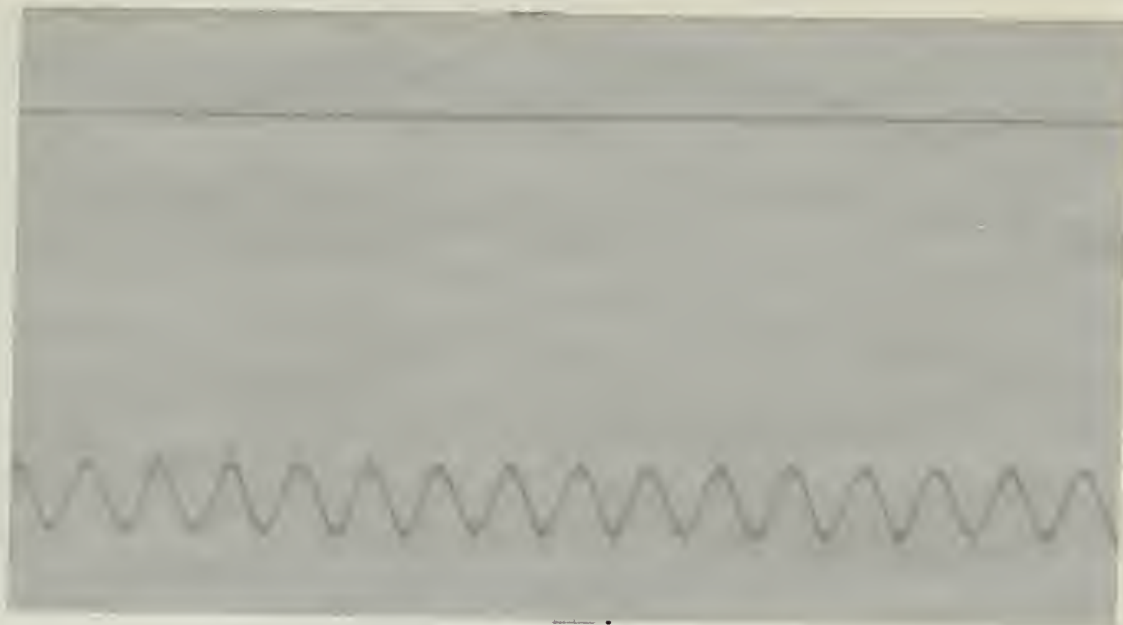
BEAM 2 - STATIC
· 156 GRAMS.

44

note: White oscillographs taken with camera of higher speed than
Camera used for black oscillographs.
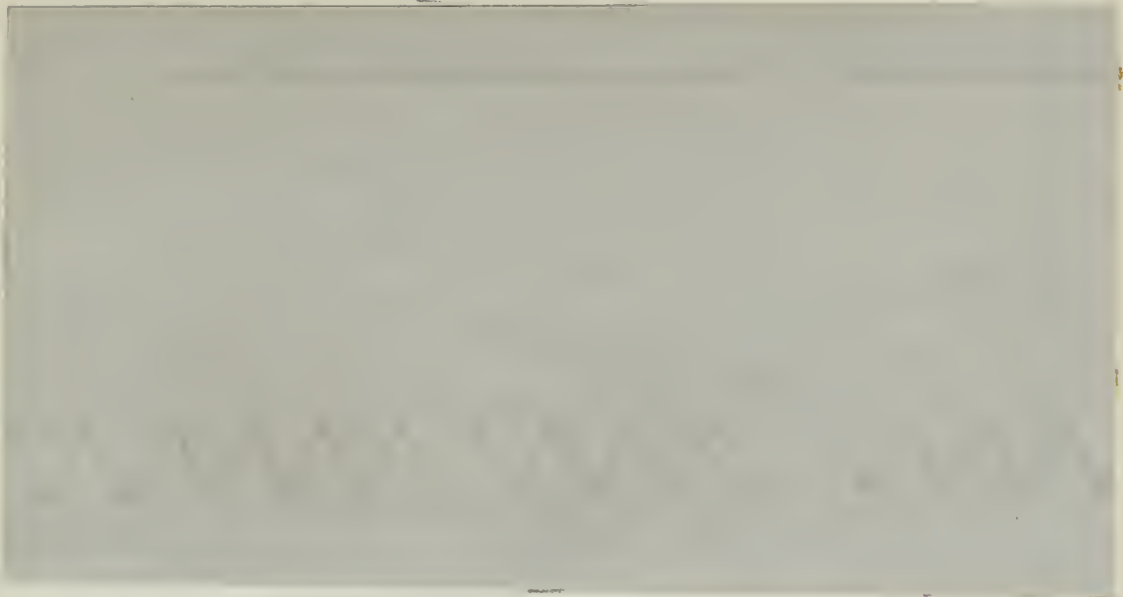
BEAM 1 - STATIC
.264 GRAMS
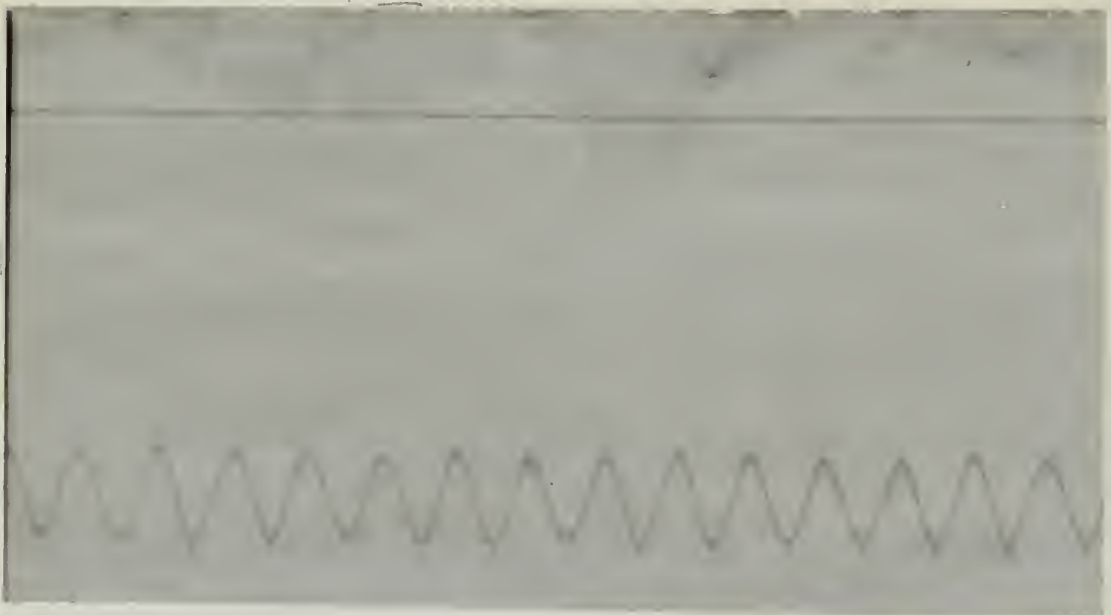
BEAM 2 - STATIC
.264 GRAMS
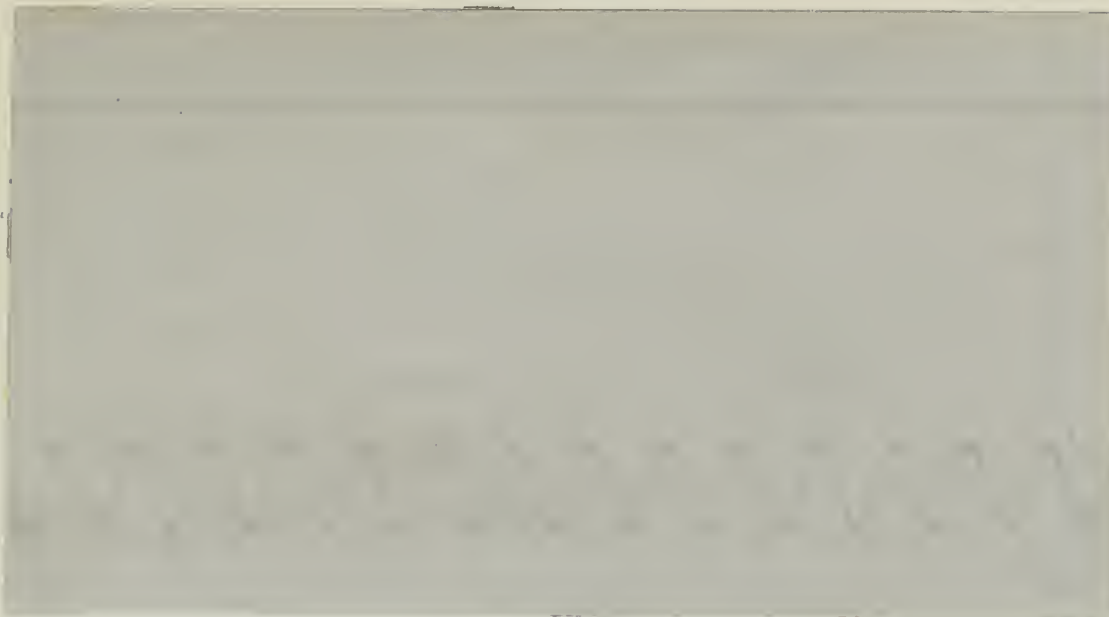
BEAM 1- STATIC
.366 GRAMS

BEAM 2 - STATIC
.366 GRAMS.
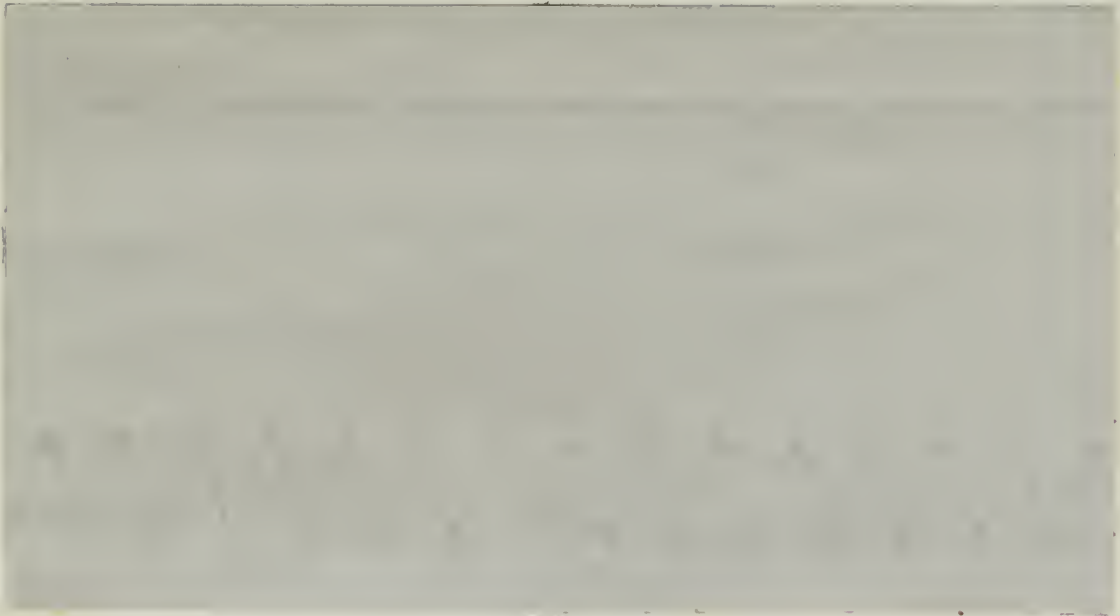
BEAM 1 - STATIC
.472 GRAMS
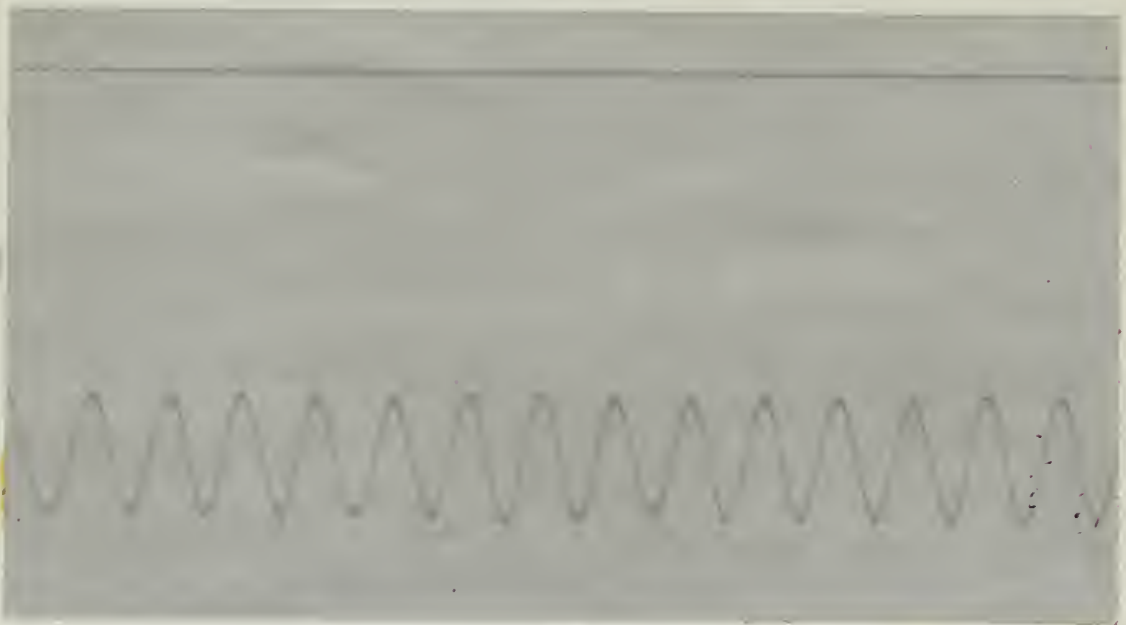
BEAM 2 - STATIC
.472 GRAMS

47
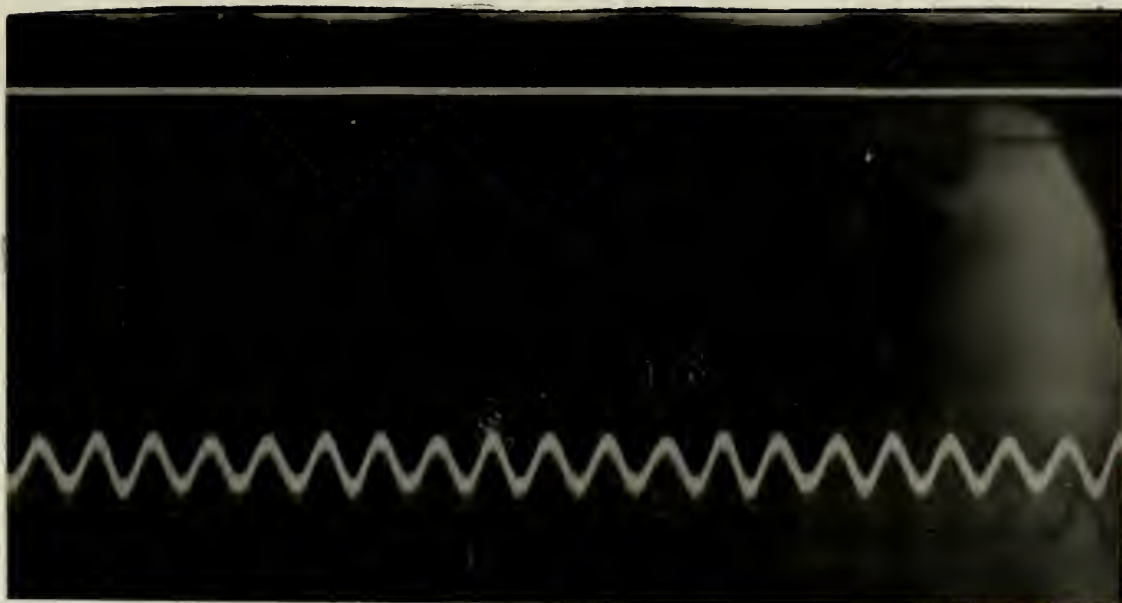
BEAM 1- DYNAMIC
· 156 GRAMS
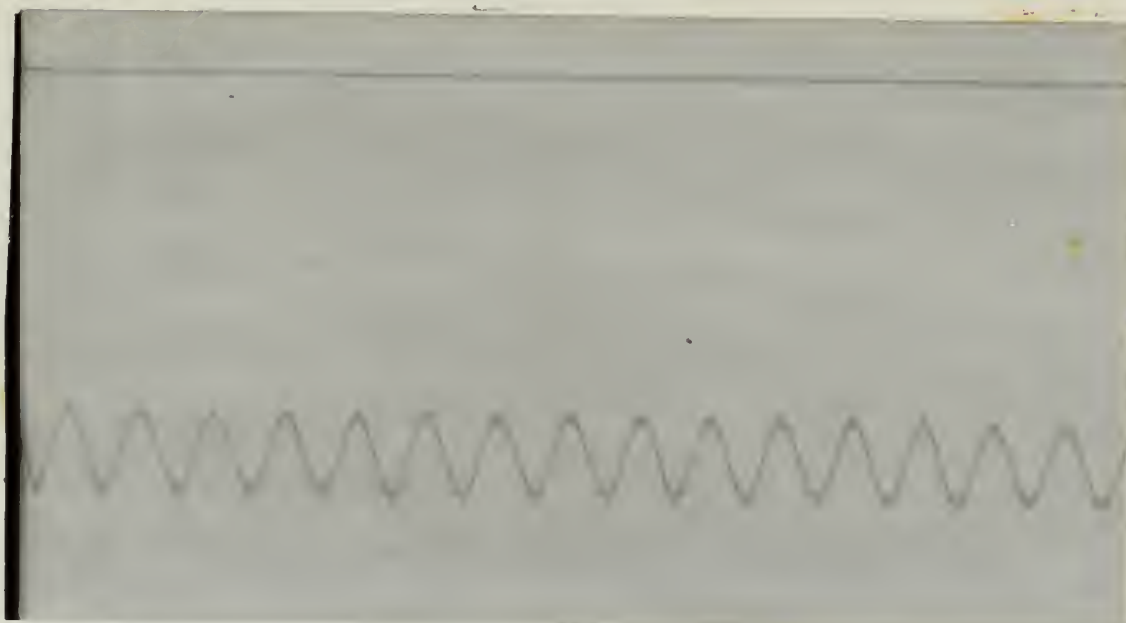
BEAM 2 - DYNAMIC
· 156 GRAMS.

48
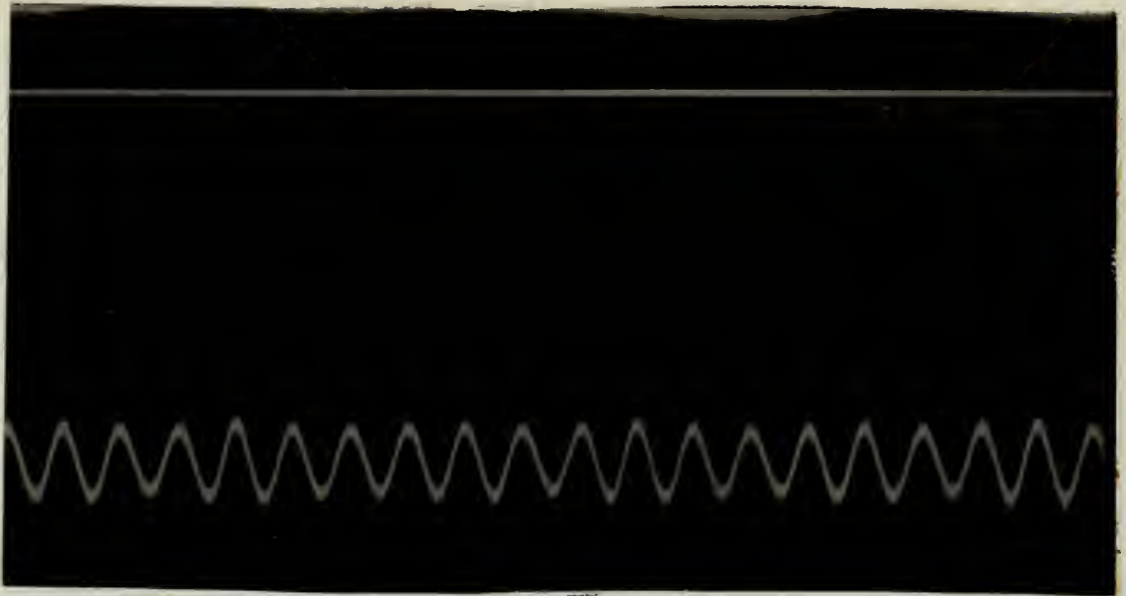
BEAM 1 - DYNAMIC
. 264 GRAMS
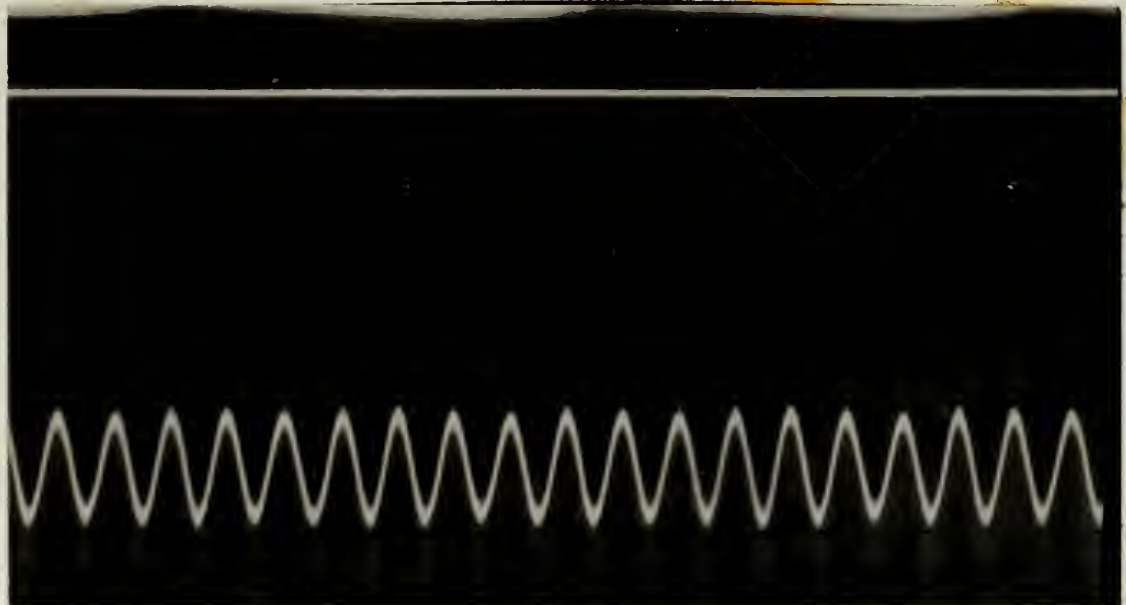
BEAM 2 - DYNAMIC
. 264 GRAMS

49

BEAM 1 - DYNAMIC
·366 GRAMS

FOR This run - extraneous
Vibrations occured. Probably
due to bearings.
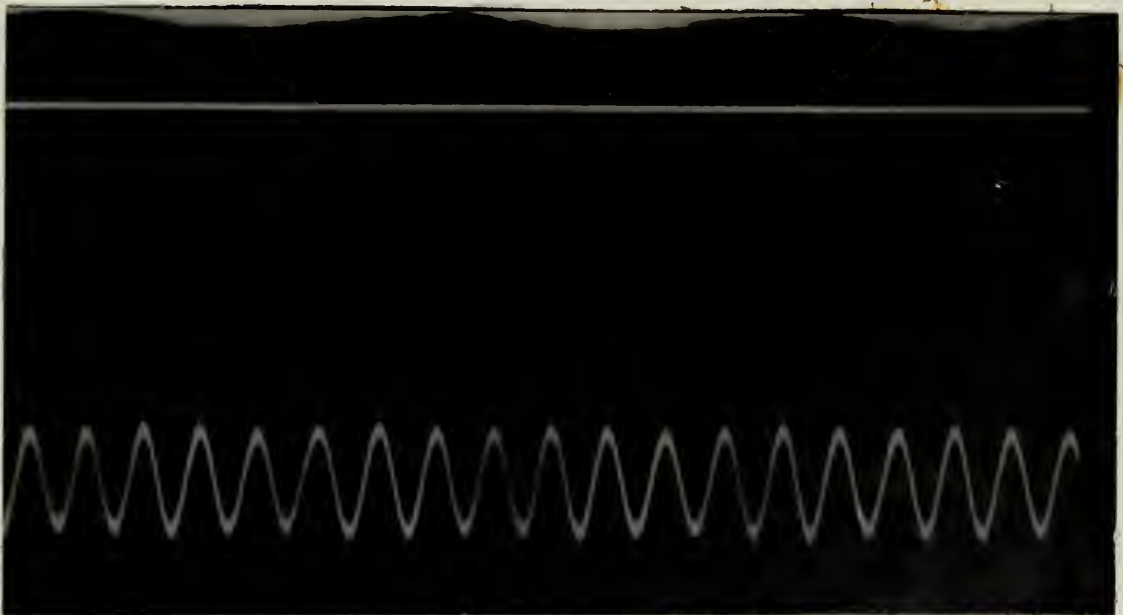
BEAM 2 - DYNAMIC
·366 GRAMS.

50
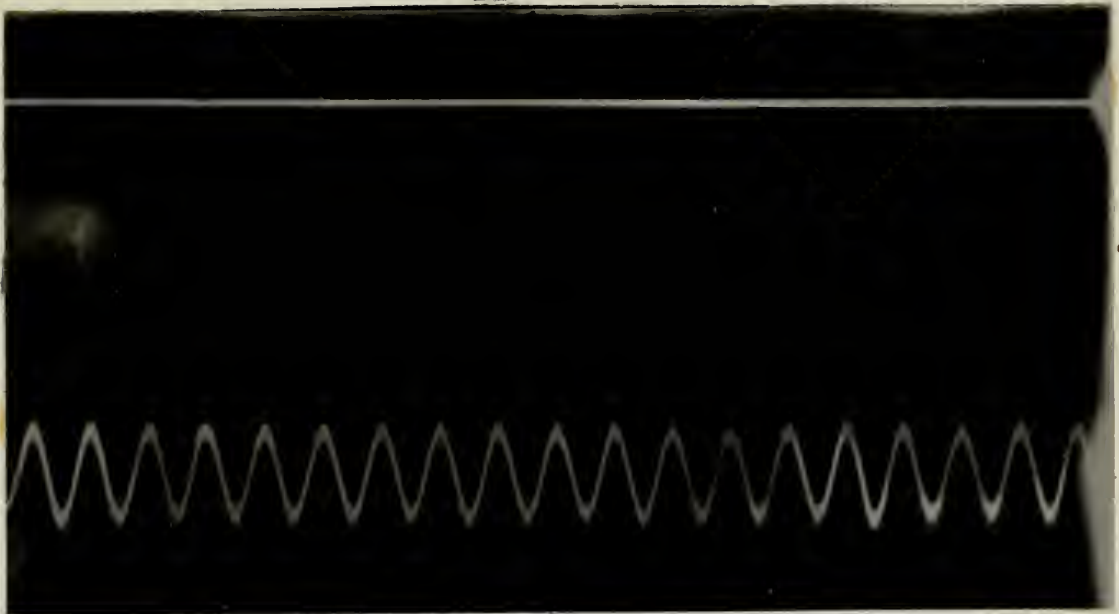
BEAM 1 - DYNAMIC
· 472 GRAMS
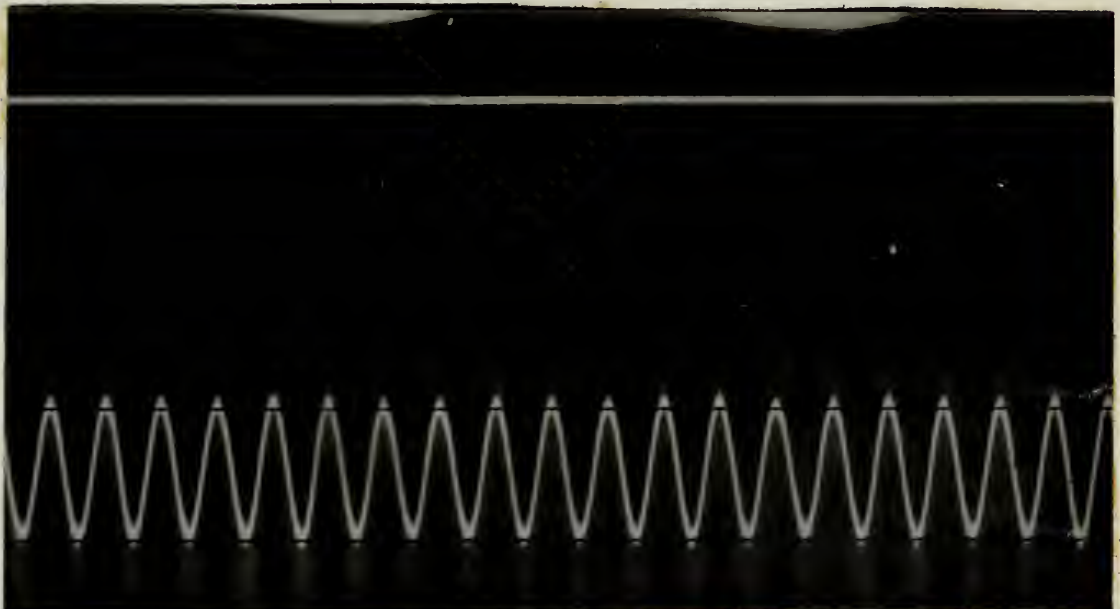
BEAM 2 - DYNAMIC
· 472 GRAMS.

51

CHAPTER VII

CONCLUSIONS AND RECOMMENDATIONS

The curves$_\wedge$ both static and dynamic unbalance indicate the
machine is quite sensitive. The smallest unbalance detected was
.594 *gm-cm* . but it is believed that if the test shaft had been
more perfectly balanced smaller unbalances could be detected.

The machine is in the stage of completion such that a rotor
with unknown unbalance could be balanced very nearly perfect.
The method of procedure would of course be by trial and error
which is time consuming. It necessitates the addition of various
size weights at different angular locations in two predetermined
planes. The rotor will be completely balanced when no wave form
and therefore no vibration is recorded. The calibration curves
would serve to give the approximate magnitude of unbalance but
would not predict the location along the axis.

To make the method more automatic, a procedure must be
devised to separate the vibrations due to static unbalance from
those due to dynamic unbalance. This would consist of an
electrical network capable of adding and subtracting impulses.
The author did not investigate this phase of the project, but
further work by members of the Mechanical Engineering Department
is contemplated.

During the course of the experiment several faults were
noticed which had a detrimental effect on the results. Most
of these can be contributed to faulty design. The results show

that one beam has a 5% greater spring constant than the other.
This difference shows up in the slopes at the calibration
curves. The difference can be contributed to variation in
dimensions and hardness of the material.

After prolonged vibration slight misalignment of the
wheels occurred which prevented smooth rotation of the rotor,
and slight clearances developed where the relatively soft aluminum
was fitted into steel.

For optimum results, steel bushings should be sweated on
the aluminum members which fit snugly into the steel base and on
the aluminum wheel shafts, which project into the inside race
of the ball bearings.

Although very small precision ball bearings were used in
a preloaded condition, excess clearance was still prevalent.
Timken tapered bearings would eliminate this to some extent.
To eliminate extraneous vibrations due to bearings, no matter
what type used, would necessitate eliminating the bearings
completely.

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                   Resonnance." Engineering, May, 1946 - Vol.161.

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                                       Flutter"

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7. Timoshenko, S. - "Vibration Problems in Engineering" -2nd edition.

Derivation for the Potential energy of a cantilever beam.

Consider the figure shown *below.*
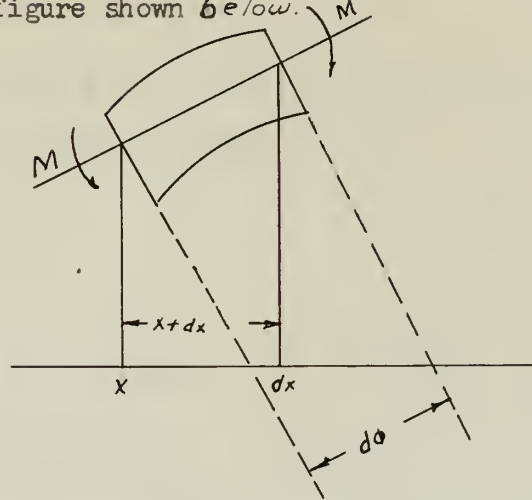
in which a small element of the beam is under the influence

of the bending moment M. The element is originally straight and

is bent through an angle $d\phi$ by the moment M. If the left-hand

end of the element is assumed to be clamped, the moment M at

the right hand end turns through the angle $d\phi$. The work done

by M on the beam is 1/2 $M\,d\phi$ where the factor 1/2 appears

because M and $d\phi$ are increasing from zero together.

This work is stored as potential energy in the beam element.

Now calculate the angle $d\phi$. If the slope at the left-

hand end x be $\dfrac{dy}{dx}$, then the slope at the right-hand end is

$\dfrac{dy}{dx} + \left(\dfrac{d^2y}{dx^2}\right) dx$ and the difference in slope $d\phi$ is

$$d\phi = \frac{d^2y}{dx^2} \, dx$$

so that $d\,pot = \frac{1}{2} M\,y'' \, dx.$

With the differential equation of bending $M = EIy''$

then $\quad Pot. = \dfrac{EI}{2} \displaystyle\int_0^\ell \left(\dfrac{d^2y}{dx^2}\right)^2 dx$

The kinetic energy for a cantilever beam with a concentrated

mass can be derived as follows:
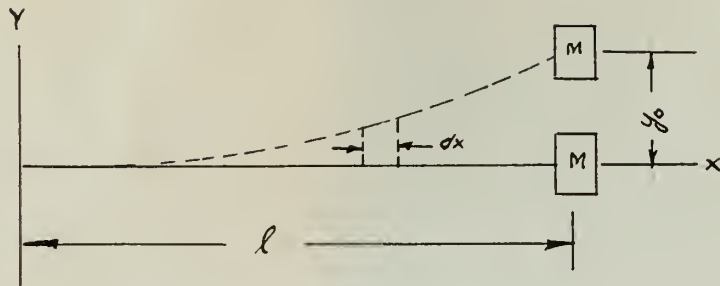
Consider an element (dx) neglecting the concentrated mass.

The Kinetic energy = $\frac{1}{2} m \omega^2$

For an elemented section KE = $\frac{1}{2} (u_1 dx)(y\omega)^2$

and the total Energy = $\frac{1}{2} u_1 \omega^2 \displaystyle\int_0^\ell y^2 dx$

Now consider a massless beam loaded with a concentrated mass.

The Kinetic Energy = $\frac{1}{2} M y_0^2 \omega^2$

Kinetic energy of system = $\frac{1}{2} u_1 \omega^2 \displaystyle\int_0^\ell y^2 dx + \frac{1}{2} M y_0^2 \omega^2$

$\frac{1}{8}D$ DRILL & TAP 36 HOLES
EQUALLY SPACED

NOTE- FINISH ALL OVER

WELD DISKS ON SHAFT.

$\frac{1}{4}$

$\frac{1}{4}$

300

300

$\frac{3}{4}$

$7\frac{1}{2}$

HARDEN & GRIND CONTACT
SURFACE

8.650

9252
9248

½"D DRILL & TAP 36 HOLES
EQUALLY SPACED

SHAFT

300

HARDEN & GRIND CONTACT
SURFACE

4

NOTE FINISH ALL OVER
WELD DISKS ON SHAFT

.500

.500

3/8

3/4

7 ½

8.650

9.252
9.248

10 ½
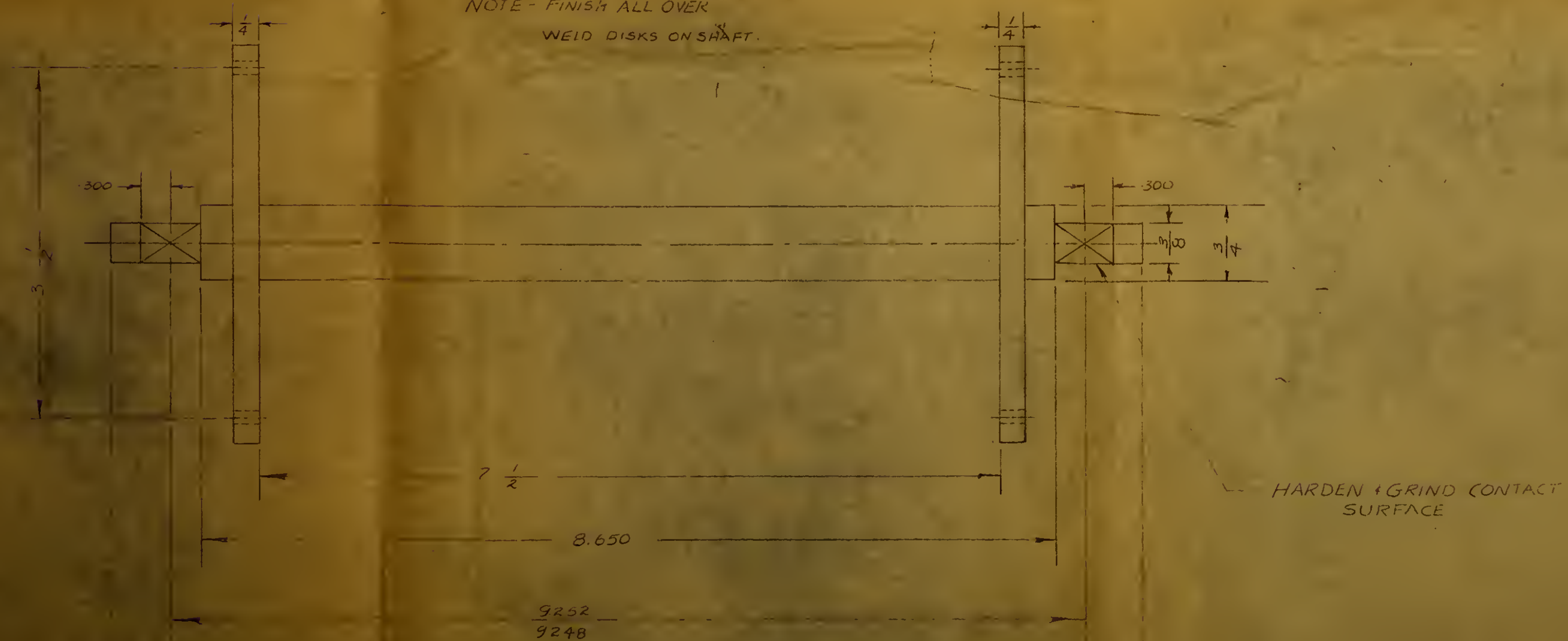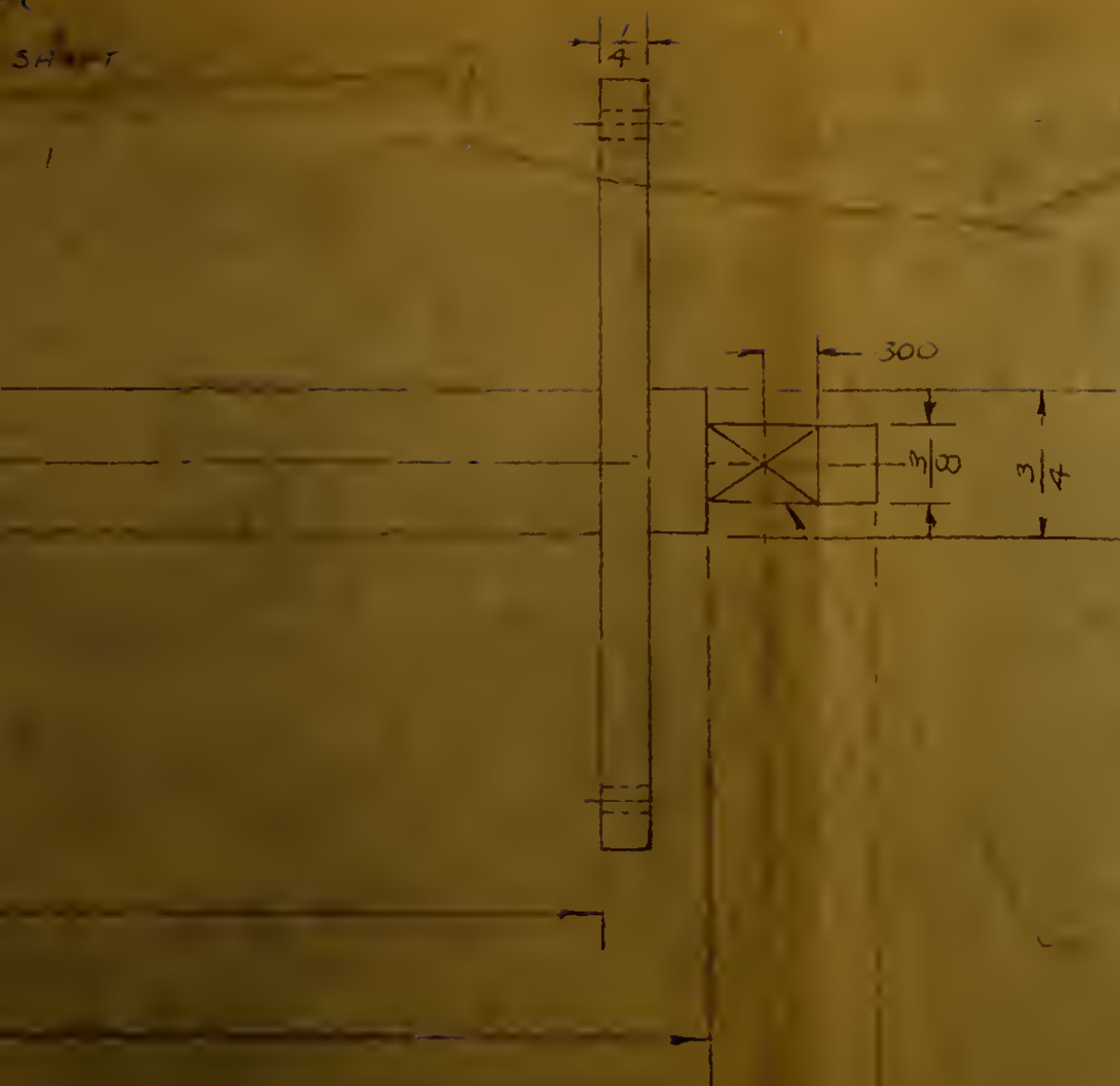
HARDEN & GRIND CONTACT
SURFACE

TEST SHAFT

HARDEN + GRIND CONTACT
SURFACE

300

.650

10 $\frac{1}{2}$

4
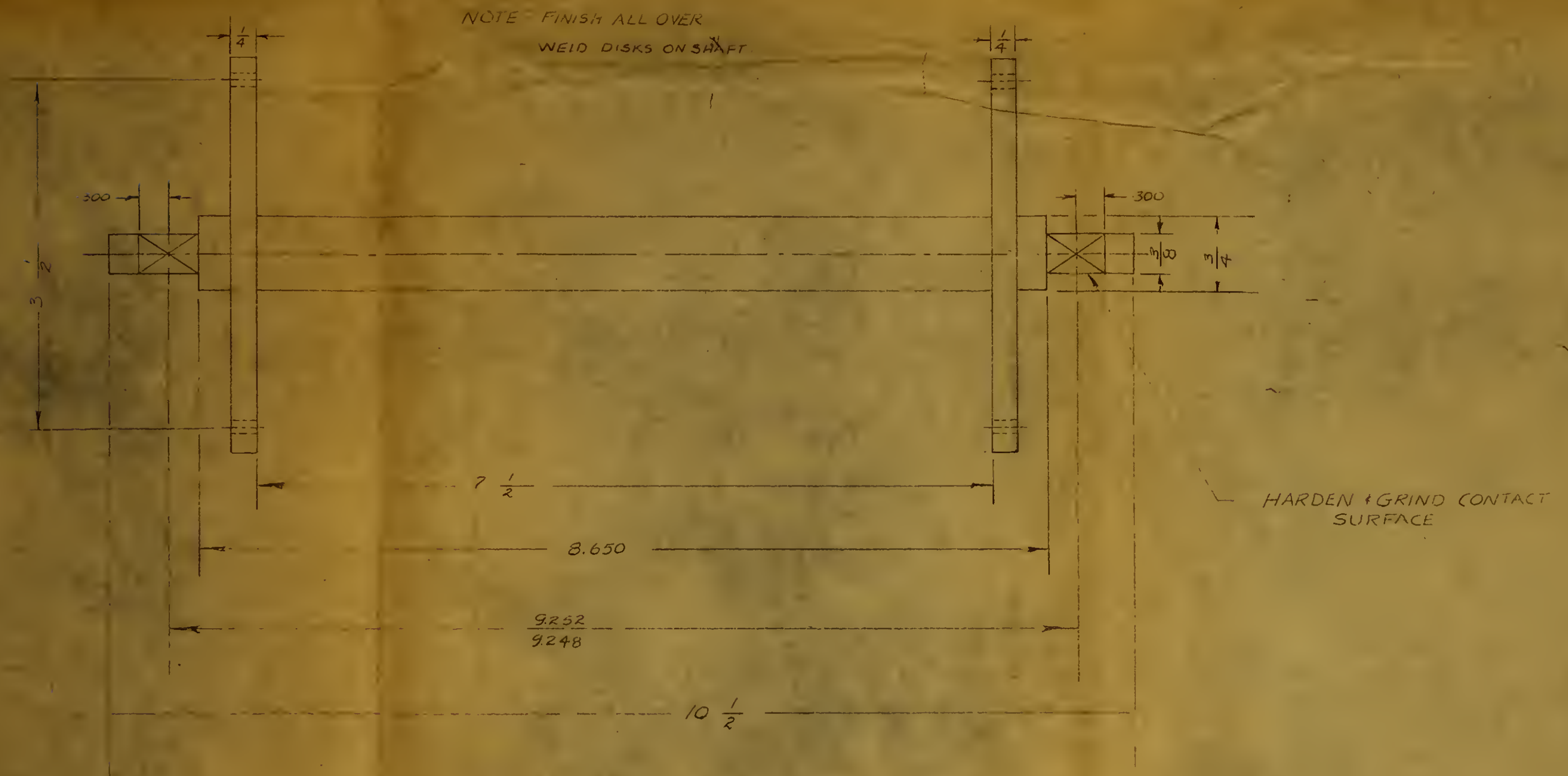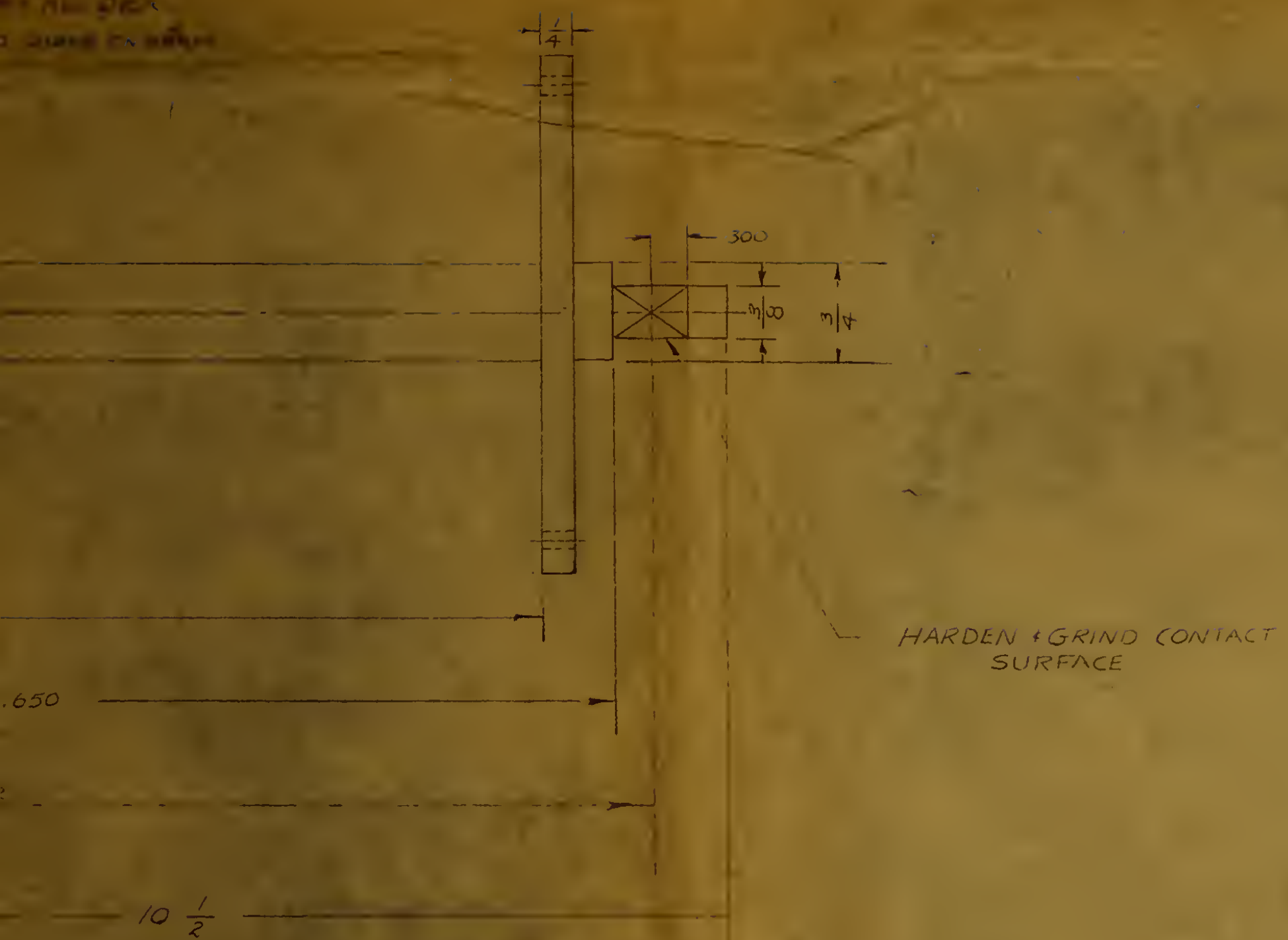
TEST SHAFT

TEST SHAFT

1 PC REQUIRED

MATERIAL - PLAIN CARBON STELL

SCALE - FULL SIZE

J.E HAMMERSTONE

‑ ‑ ‑ ½ DRILL & ASSEMBLY TO LOWELL PINS 1" DEEP

HIS 4" DEEP

.660 D DRILL THESE 6 HOLES
.600

$2\frac{3}{16}$ $\frac{9}{16}$

$\frac{G}{13}$

$2\frac{1}{2}$

$3\frac{1}{2}$

13

BACK HALF LEFT WHEEL BRACKET

FRONT HALF RIGHT WHEEL BRACKET

BACK HALF RIGHT WHEEL BRACKET

.050 WIDTH UNDERCUT OR RAD
FILLET FOR GRINDING RELIEF

$\frac{3}{32}$ D DRILL

.050 WIDTH UNDERCUT OR RAD
FILLET FOR GRINDING RELIEF

.1040 D - 36 NF2 DRILL + TAP
4 HOLES 300° DEEP
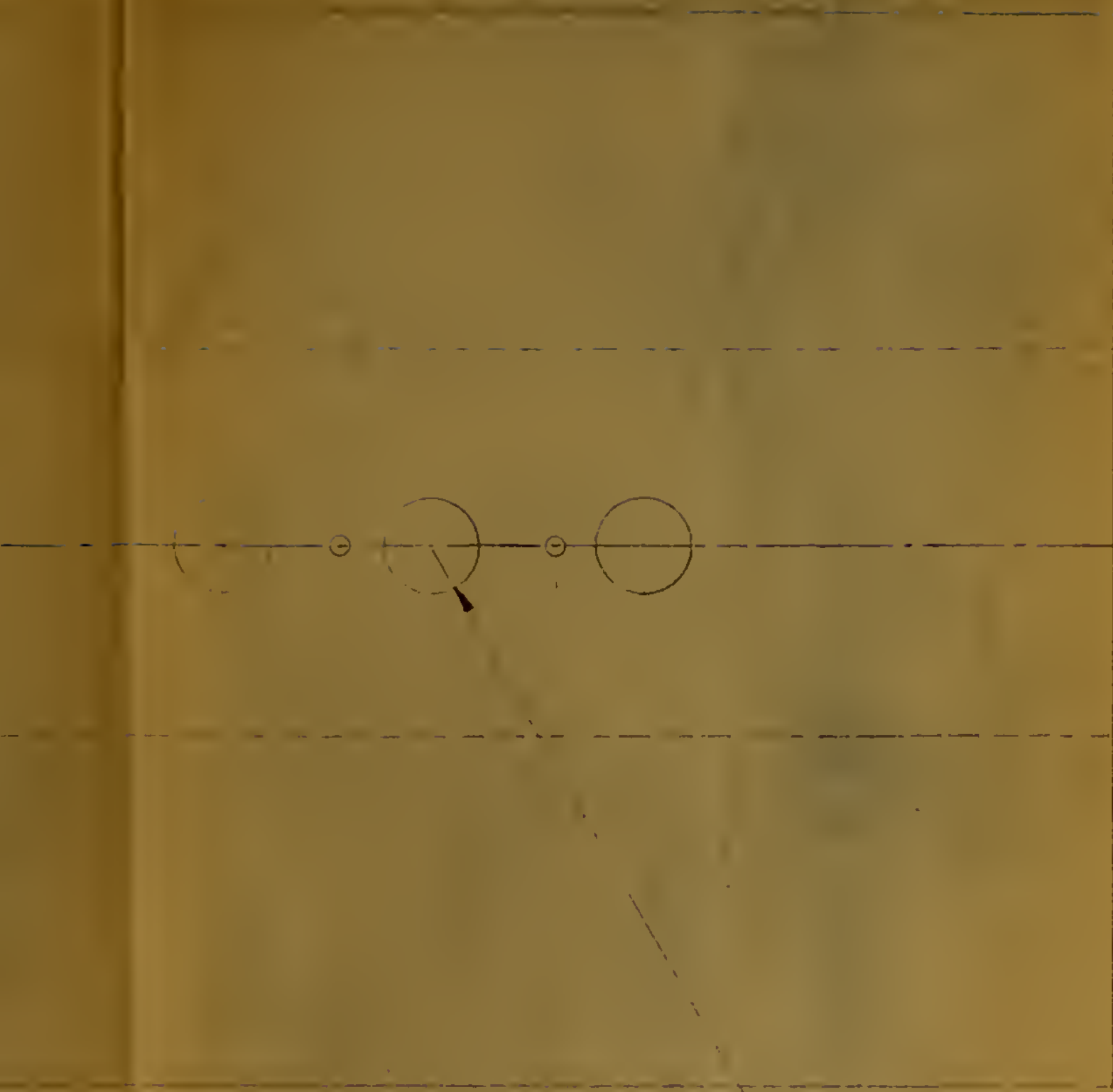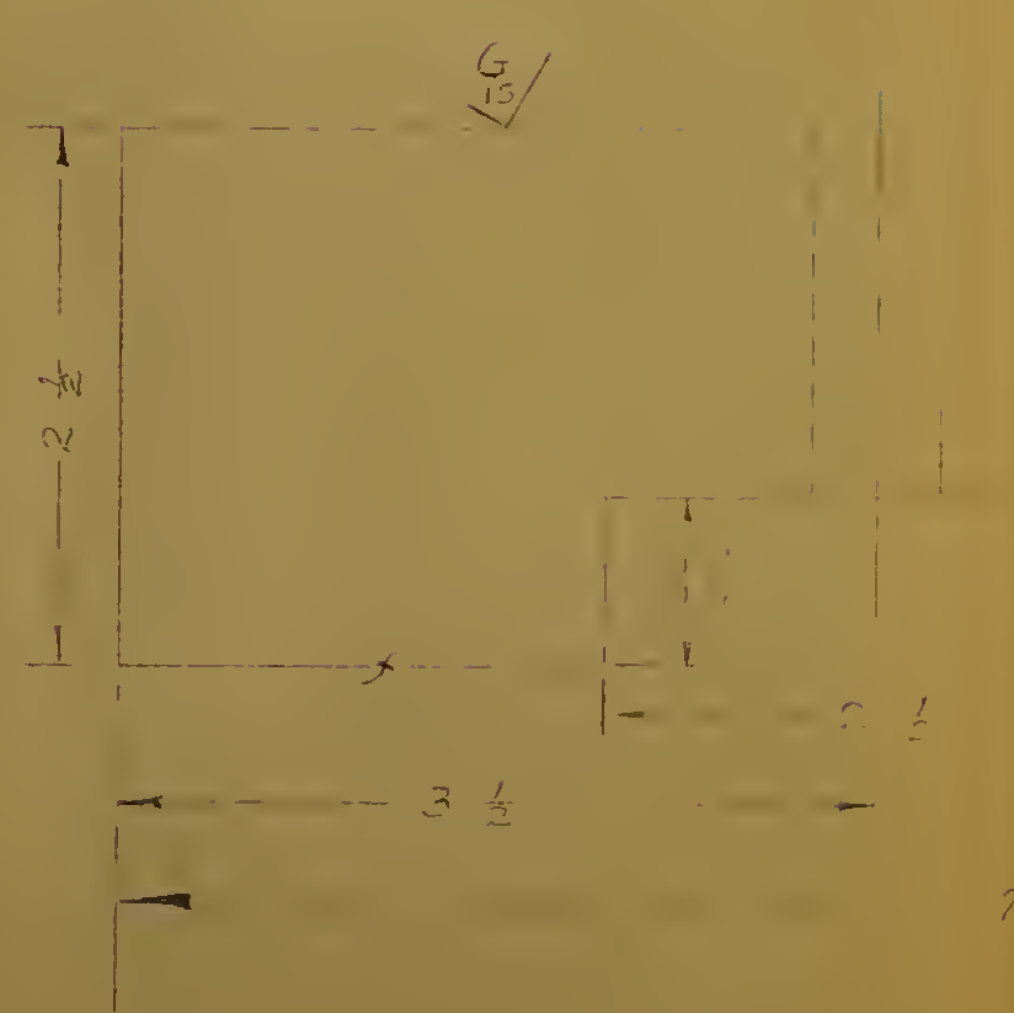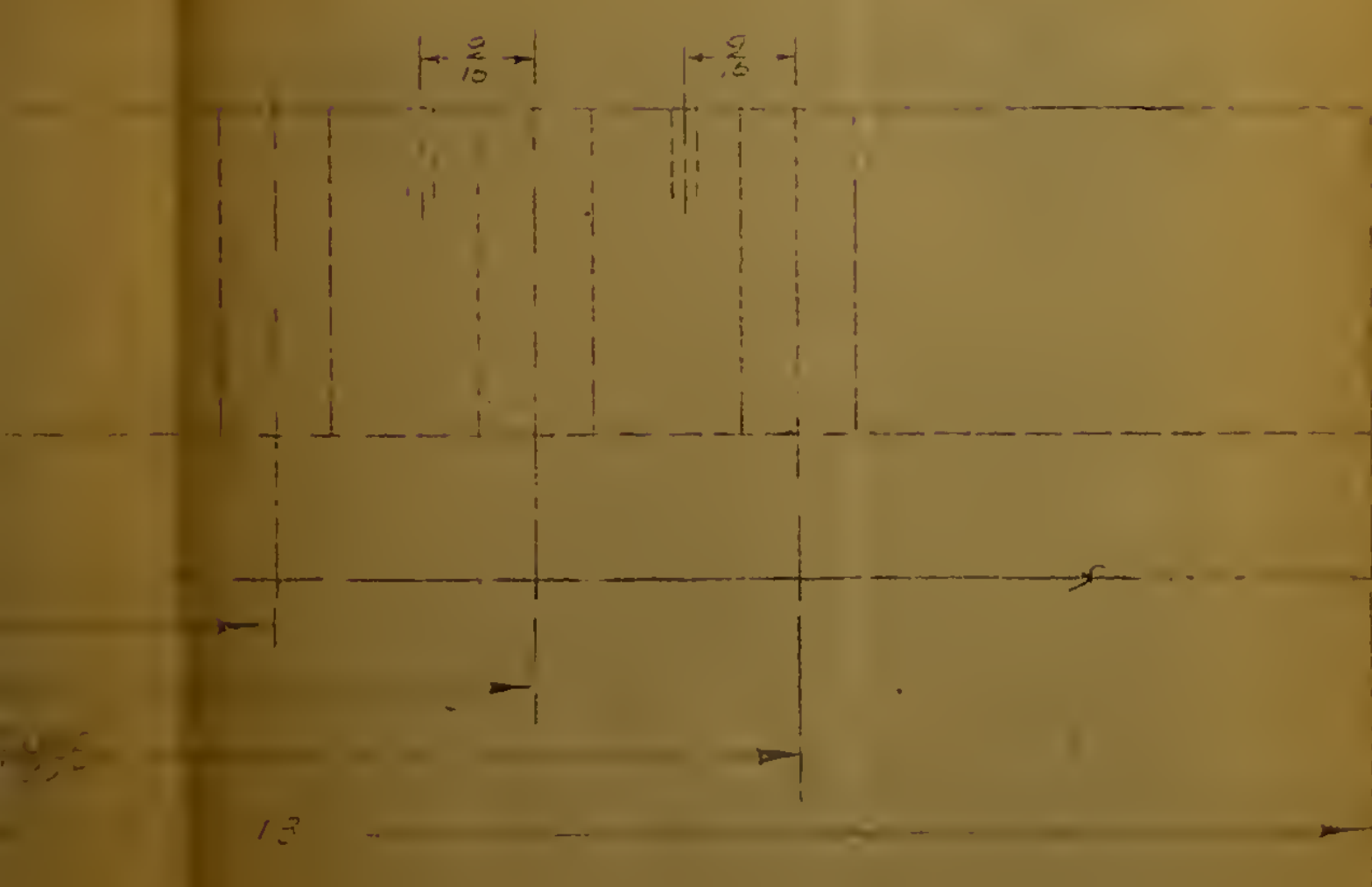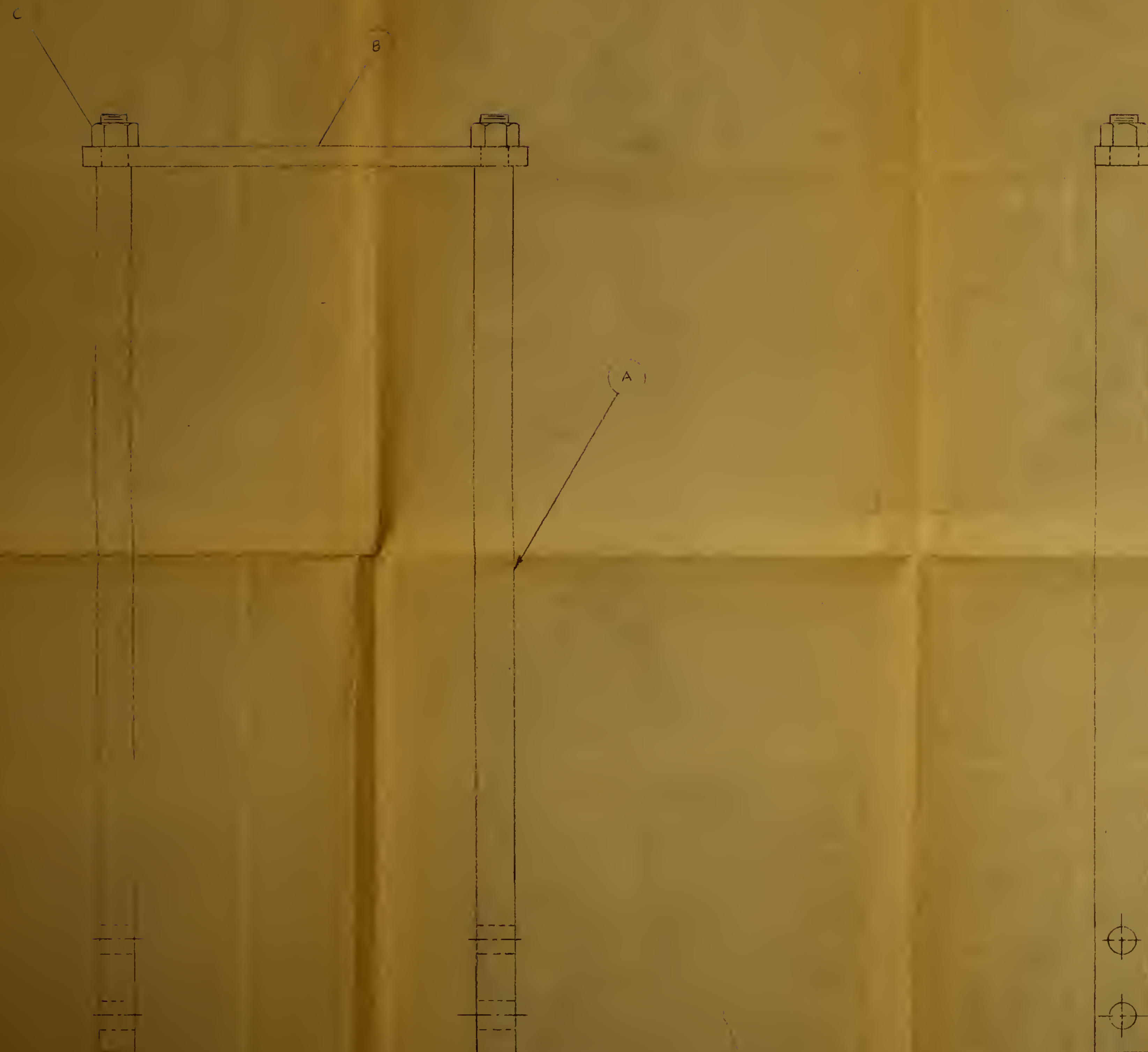
BILL OF MATERIALS			
PC NO	NAME	Q	MATERIAL
A	UPRIGHT	4	PLAIN STEEL
B	CROSS PIECE	2	PLAIN STEEL
C	3/8 A.S NUT	4	PURCHASE

J.E. HANMER 1946

CROSS SUPPORT

$\frac{0030}{0047}$

$\frac{1260}{1250}$

$\frac{0672}{1640}$ DRILL / REAM

$\frac{1993}{1990}$ $\frac{4005}{4000}$ $\frac{1601}{1600}$

$\frac{2 RUN}{2.000}$

$\frac{.061}{.060}$

$\frac{501}{500}$

$\frac{501}{500}$

$\frac{.375}{.374}$

$\frac{.950}{.949}$

$\frac{.250}{.249}$

$\frac{.250}{.249}$ $\frac{.551}{.550}$

$\frac{1101}{1100}$

$\frac{1.0274}{1.0272}$
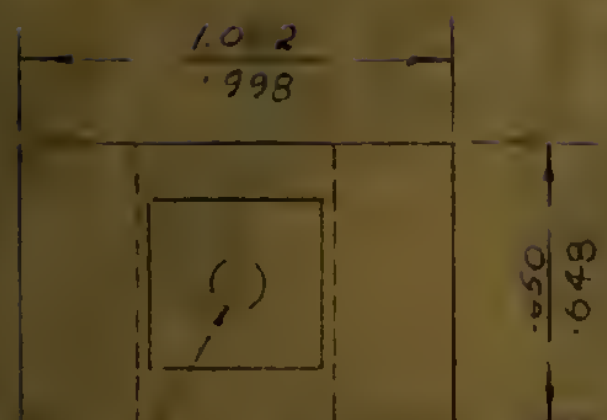
FRONT HALF LEFT WHEEL BRACKET

$\frac{1.002}{.998}$

$\frac{.650}{.648}$

LAYOUT FOR CUTTING EDGE
OF MILLING CUTTER
SCALE 4 X ACTUAL

32°

$7\frac{1}{2}$° $24\frac{3}{4}$°
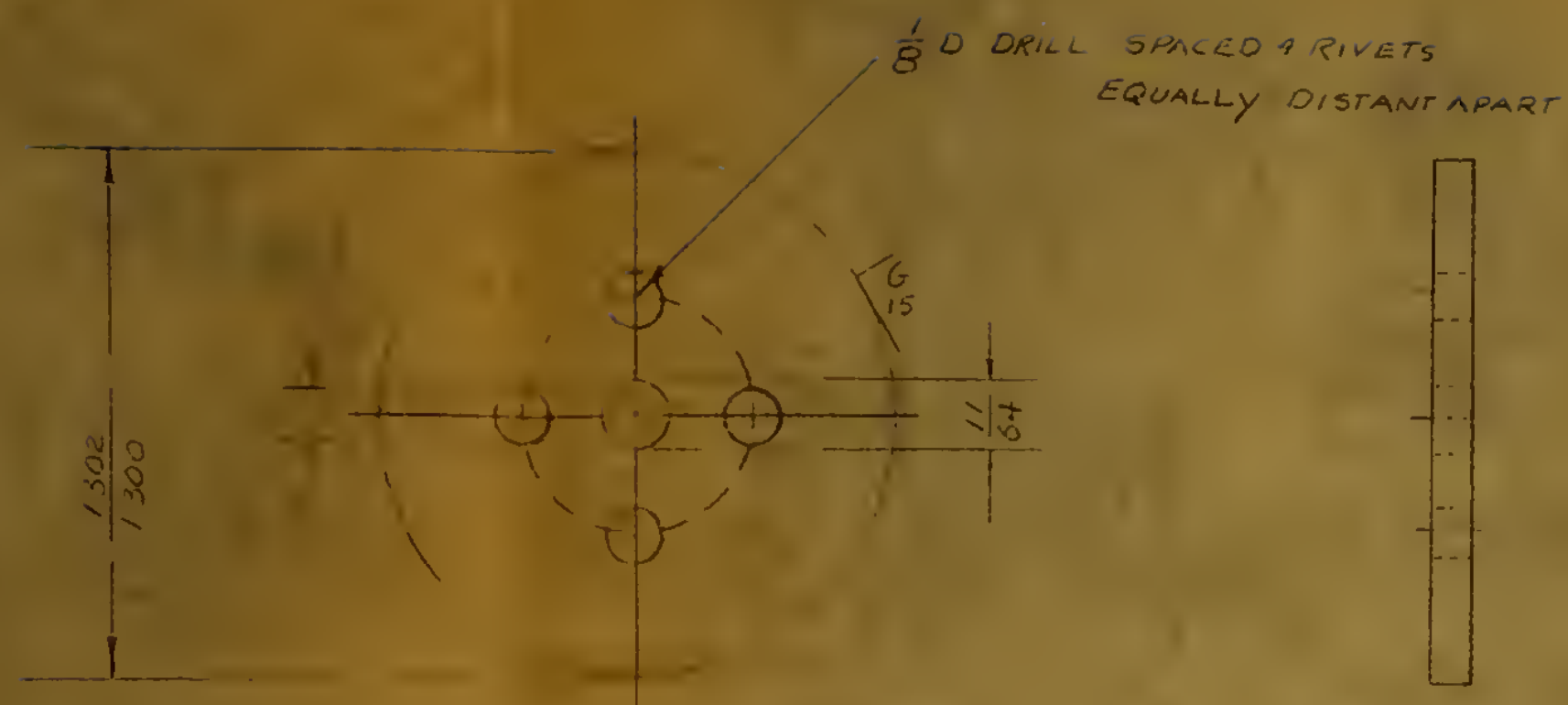
FLANGES TO BE PRESSED ON BUCKET WHEEL
2 REQUIRED PER WHEEL

GRIND .005 FROM BUCKETS
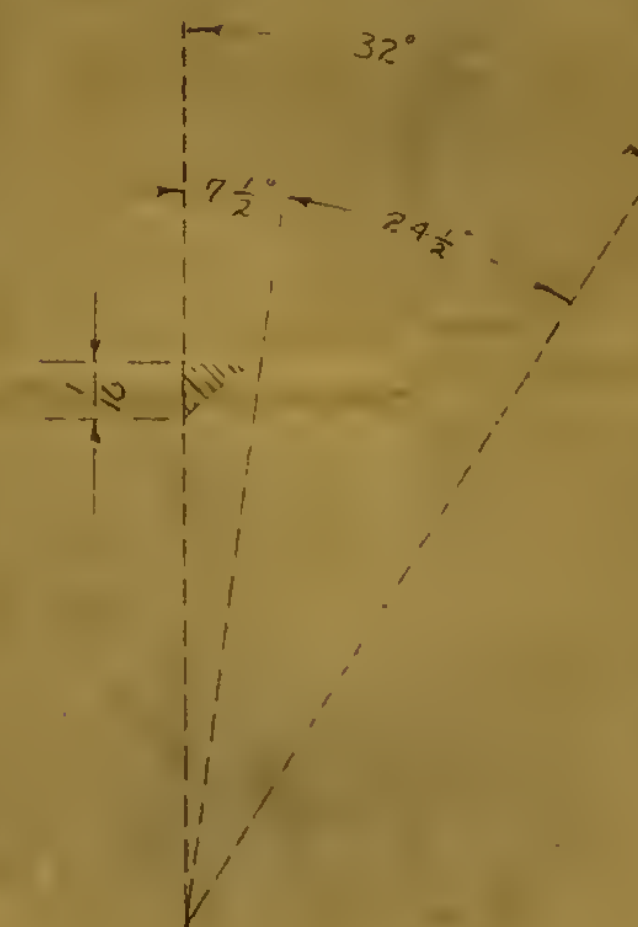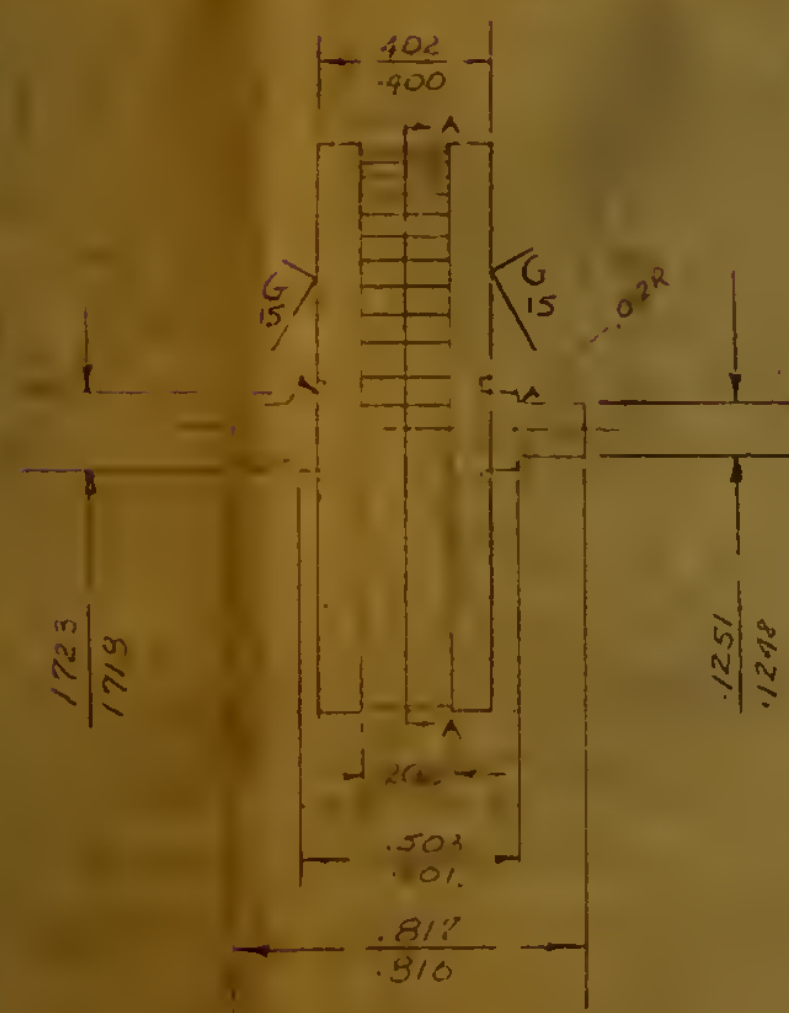
NOTE - CUT BUCKETS EVERY $7\frac{1}{2}$ DEGREES
AROUND ENTIRE PERIPHERY

.402
.400

$\frac{1}{8}$ D DRILL IN
ASSEMBLY FOR RIVETS

SECTION A·A

.1251
.1248

.1723
.1719

.504
.501

.817
.816
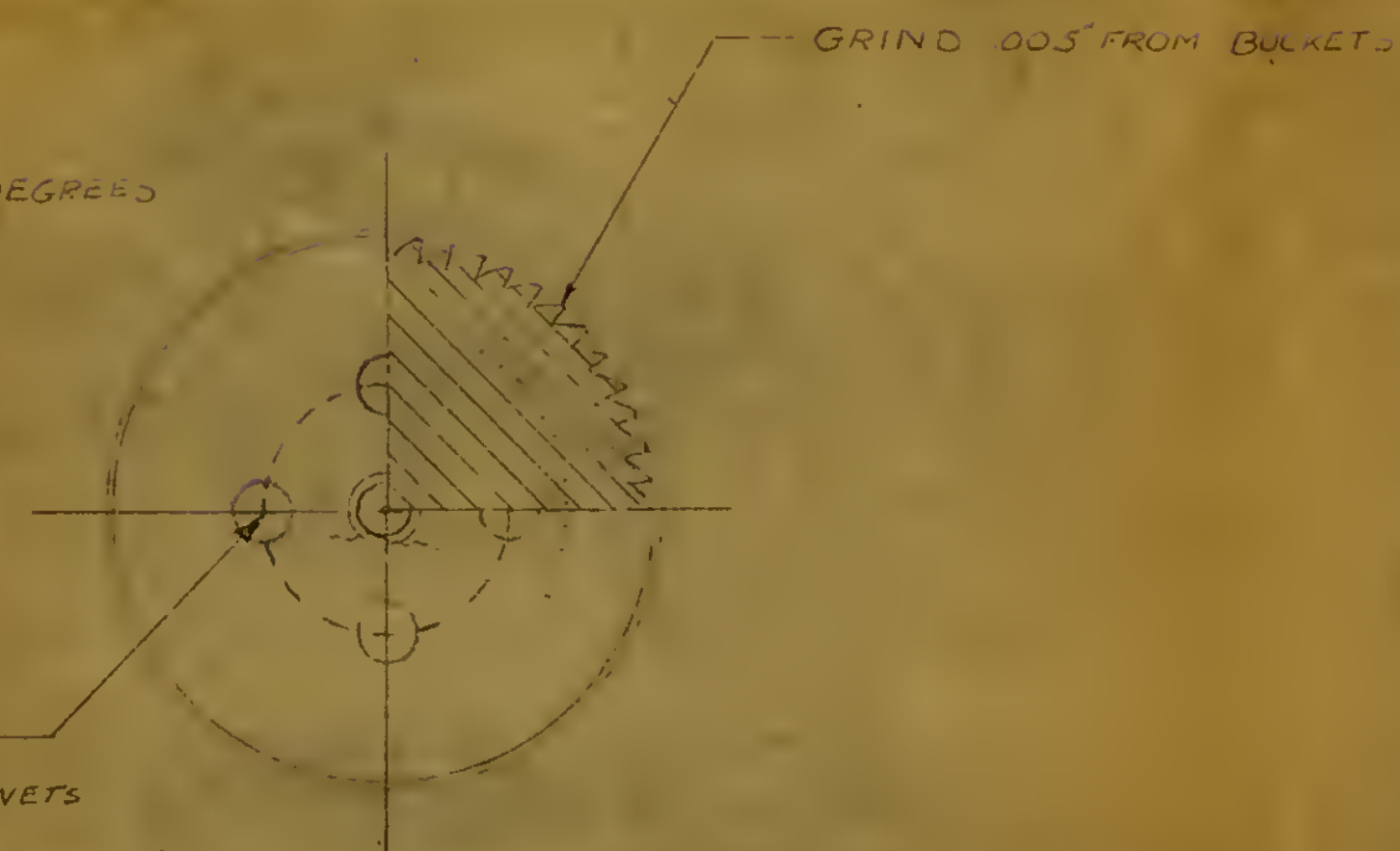
WHEEL ASSEMBLY

LAYOUT FOR CUTTING EDGE
OF MILLING CUTTER
SCALE 4 X ACTUAL

32°

7½°   24½°

GRIND .005" FROM BUCKETS

SECTION A·A

DRILL IN
ASSEMBLY FOR RIVETS

BEARING COVER

COUNTERBORE .05 DEEP

$\frac{3}{8}$ 24-NF2

$\frac{950}{949}$

$\frac{950}{948}$

$\frac{4531}{4528}$

$\frac{3}{8}$

$\frac{1}{8}$ R

$\frac{1}{8}$ D DRILL

G
15

$\frac{500}{498}$

$\frac{750}{748}$

$10\frac{3}{8}$

$\frac{6.752}{6.748}$

$\frac{7.352}{7.350}$

$\frac{7.002}{6.998}$

G
15

G
15

G
15

G
15

.050 WIDTH UNDERCUT OR RAD
FILLET FOR GRINDING RELIEF

$7\frac{1}{2}$V

.300

G
15

$2\frac{1}{2}$

G
15

$\frac{1}{2}$ D 20-NF-R

$\frac{1}{8}$ F

$\frac{1}{2}$

$\frac{5842}{5839}$

$1\frac{13}{32}$
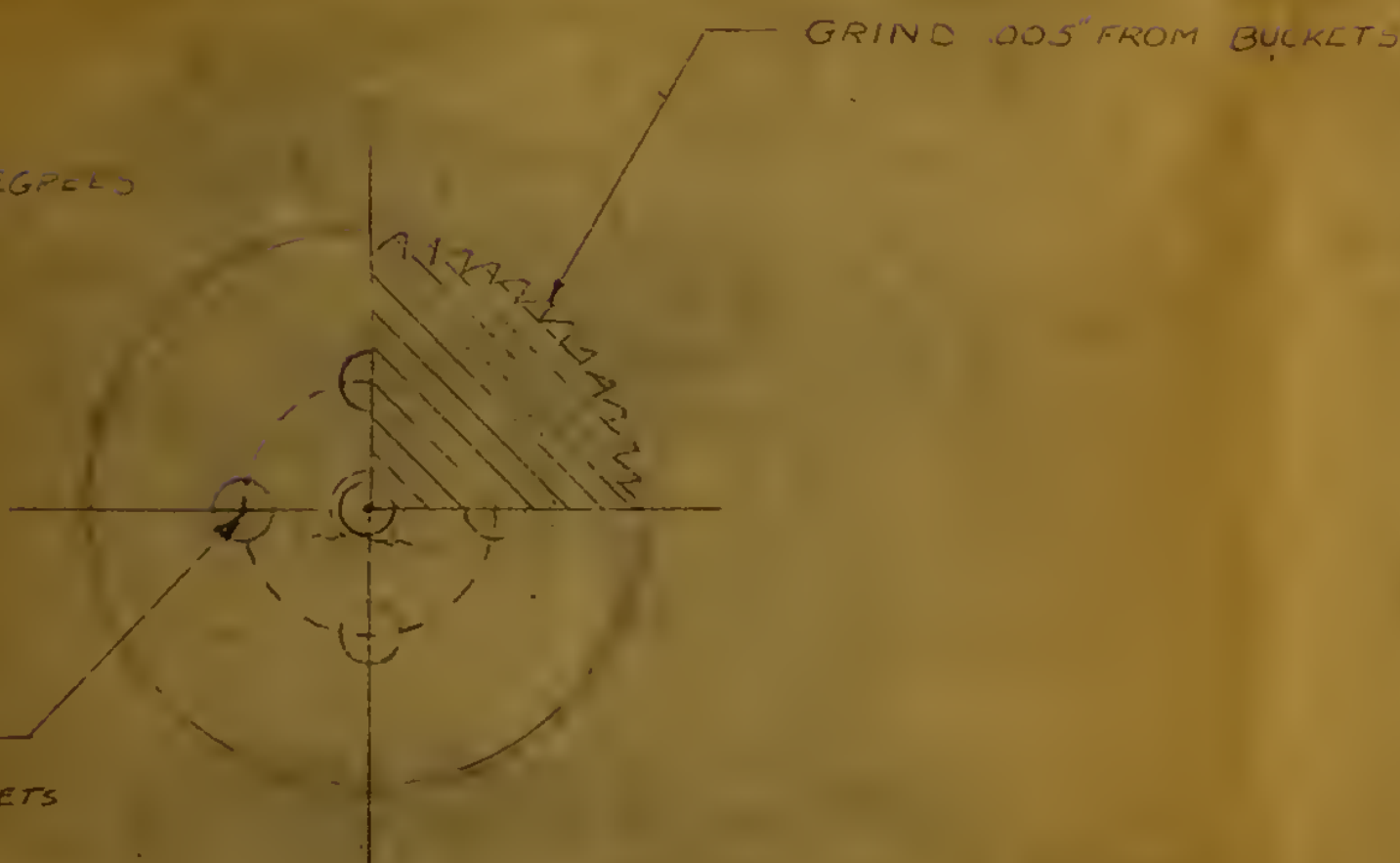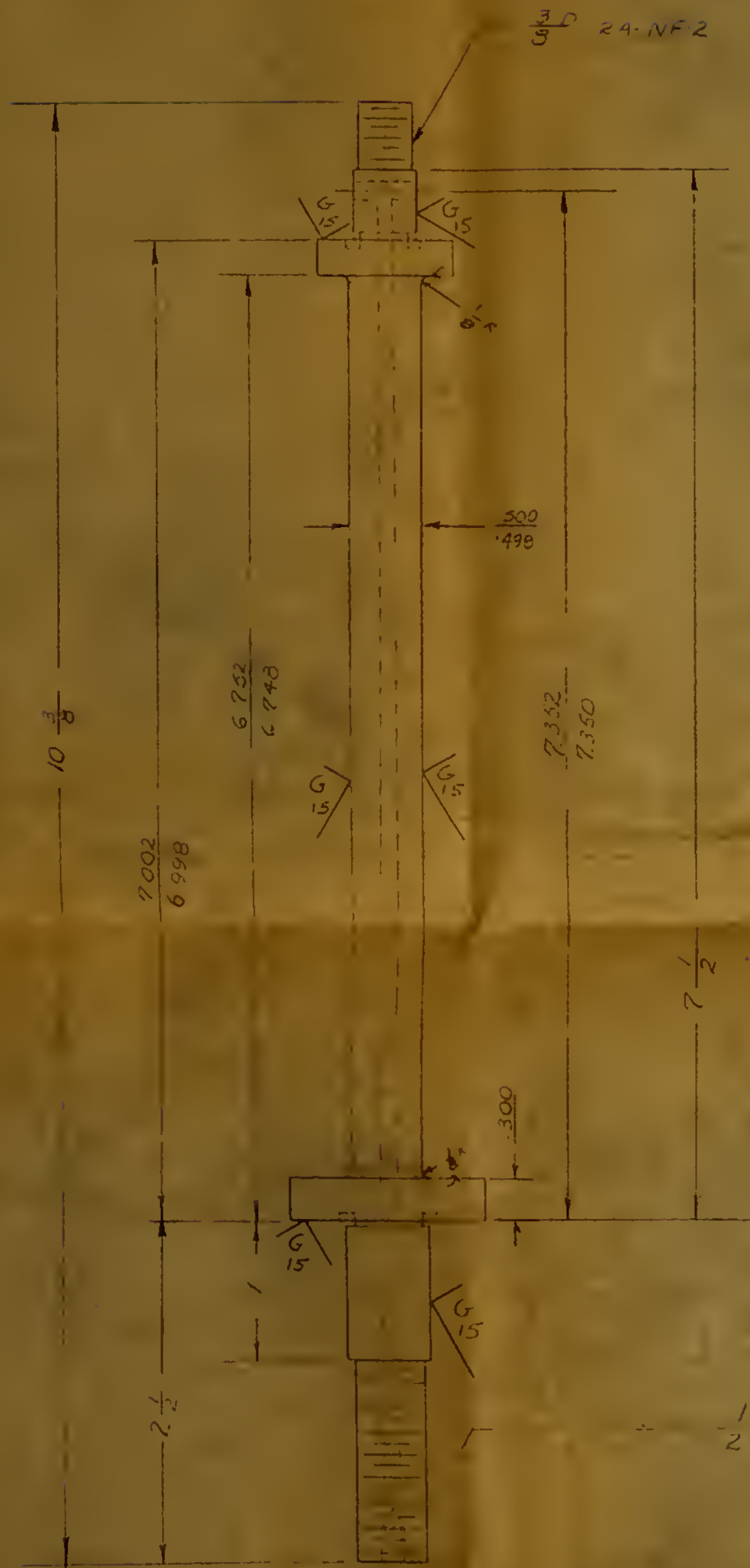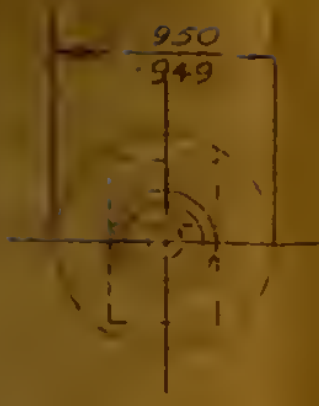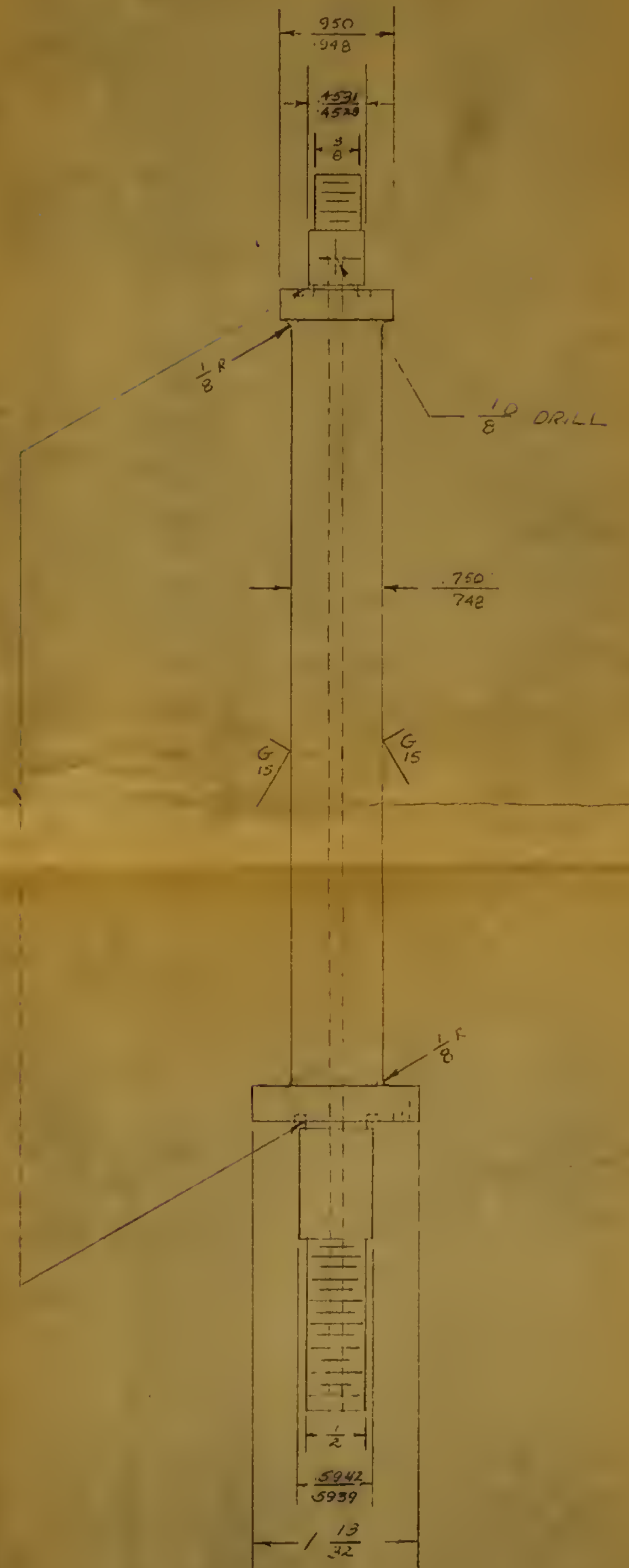
DRILL ASSEMBLY FOR DOWEL PIN
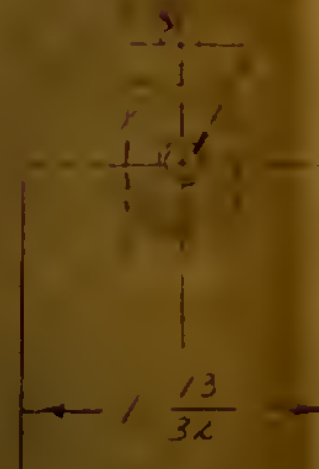
$\frac{5}{32}$ D DRILL & TAP $\frac{3}{8}$ DEEP
FOR AIR CONNECTION

$1\frac{13}{32}$

.950
.948

.4531
.4528

3/8

1/8 R

1/8" DRILL

.750
.748

DO NOT UNDERCUT OR RAI
FILLET FOR GRINDING RELIE

G 15

G 15

1/8 R

1/2

.5942
.5939

1 13/32

SUPPORT
MATERIAL
DASH REQUIRE